

DNA supercoiling: plectonemes or curls ?

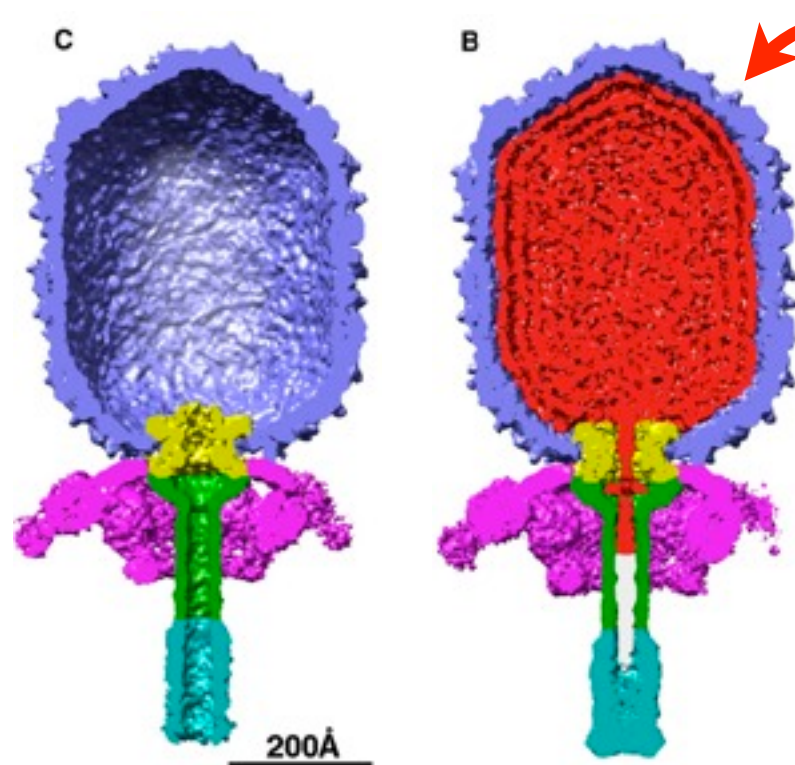
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Northwestern University (IL. USA)

Why study DNA mechanical properties ?

mechanical properties influence biology of the cell

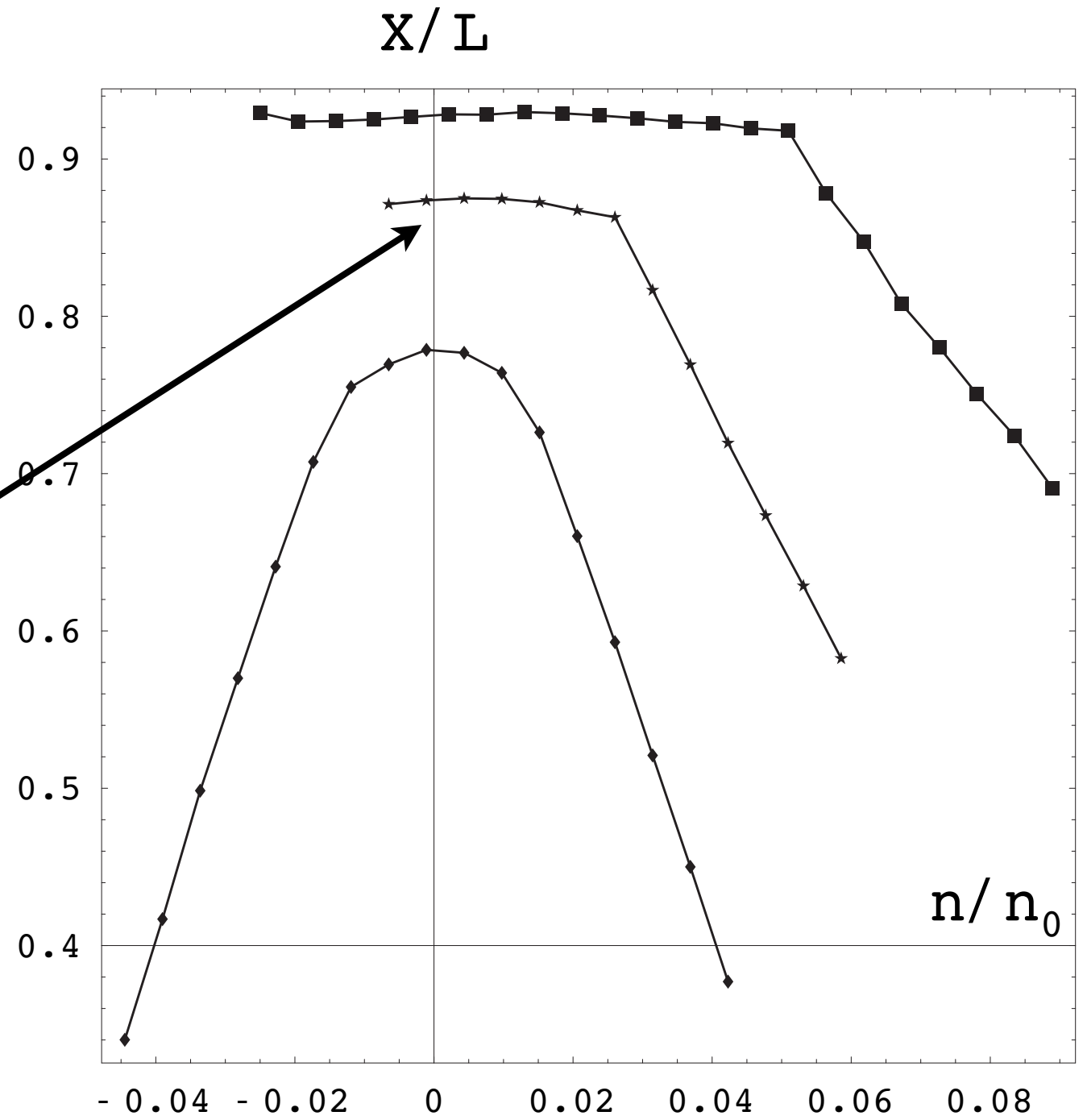
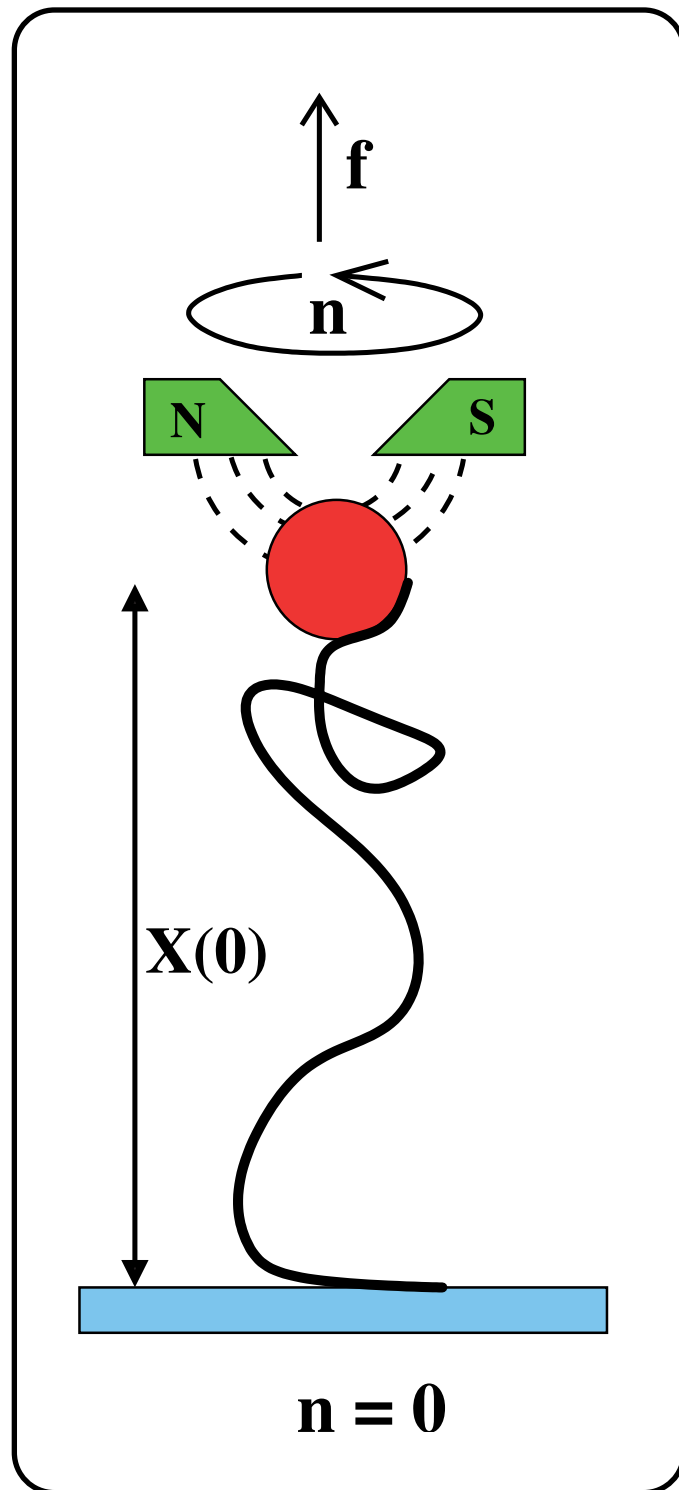
- 1 meter of DNA in a 10 micron wide nucleus
- ejection from viral capsid
- transcription (RNAPolymerase is torque dependent)
- protein binding is strain dependent, or induces strain on DNA
- chromatin compaction/decompaction (cell division)



20kbp dsDNA (6800nm)
in 40nm wide capsid

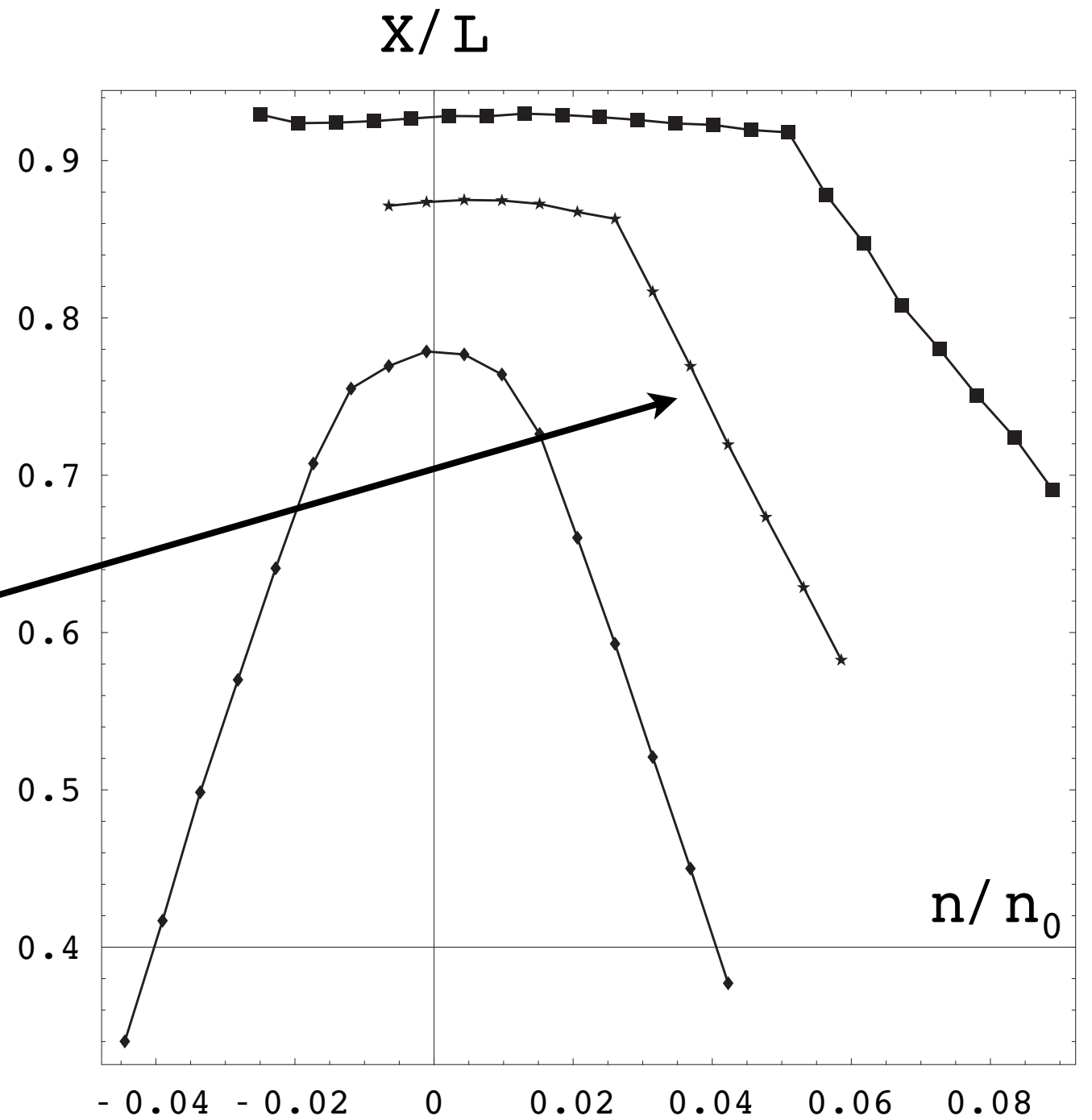
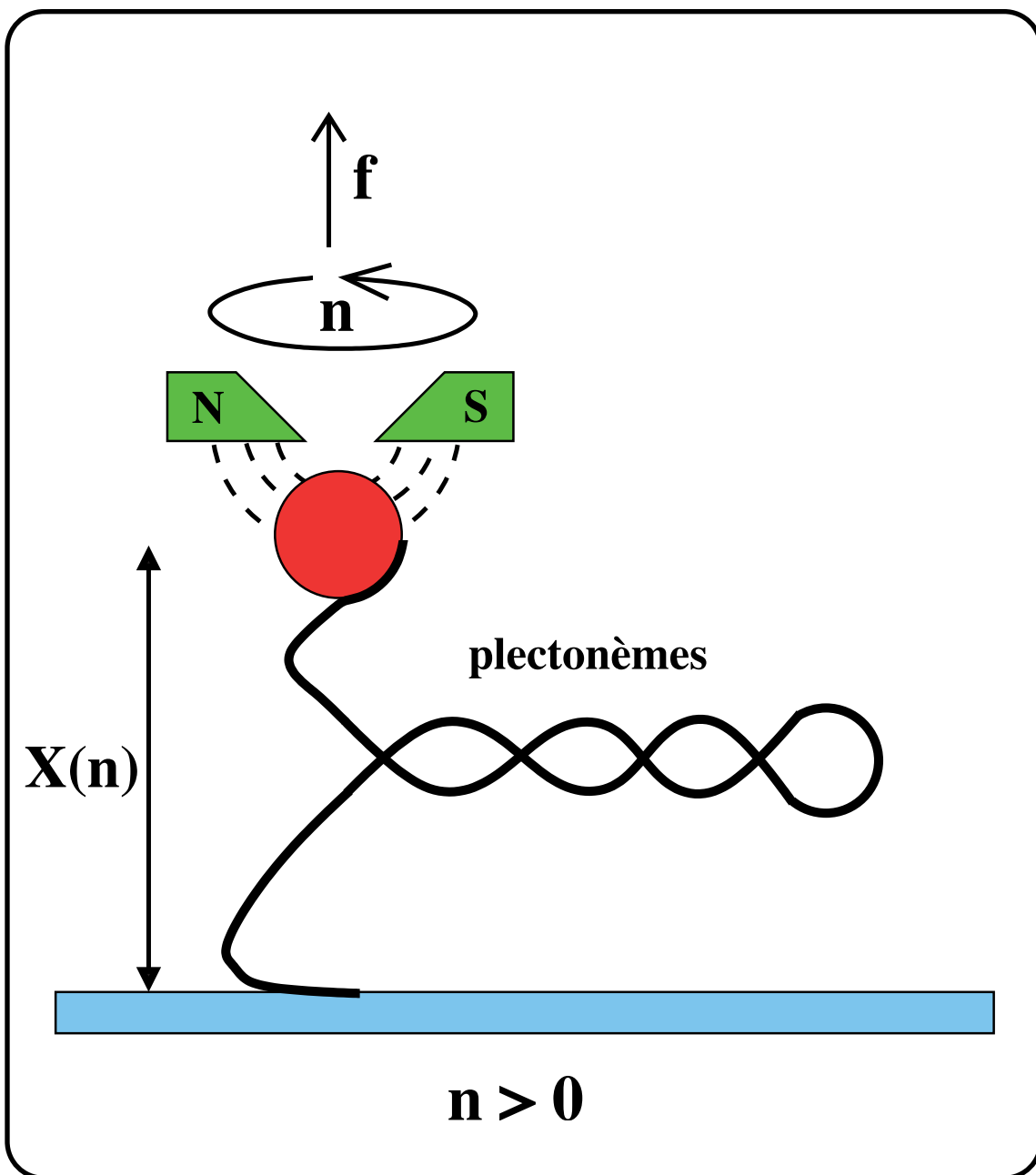
Tang et al (Structure) 2008

Pulling and twisting DNA



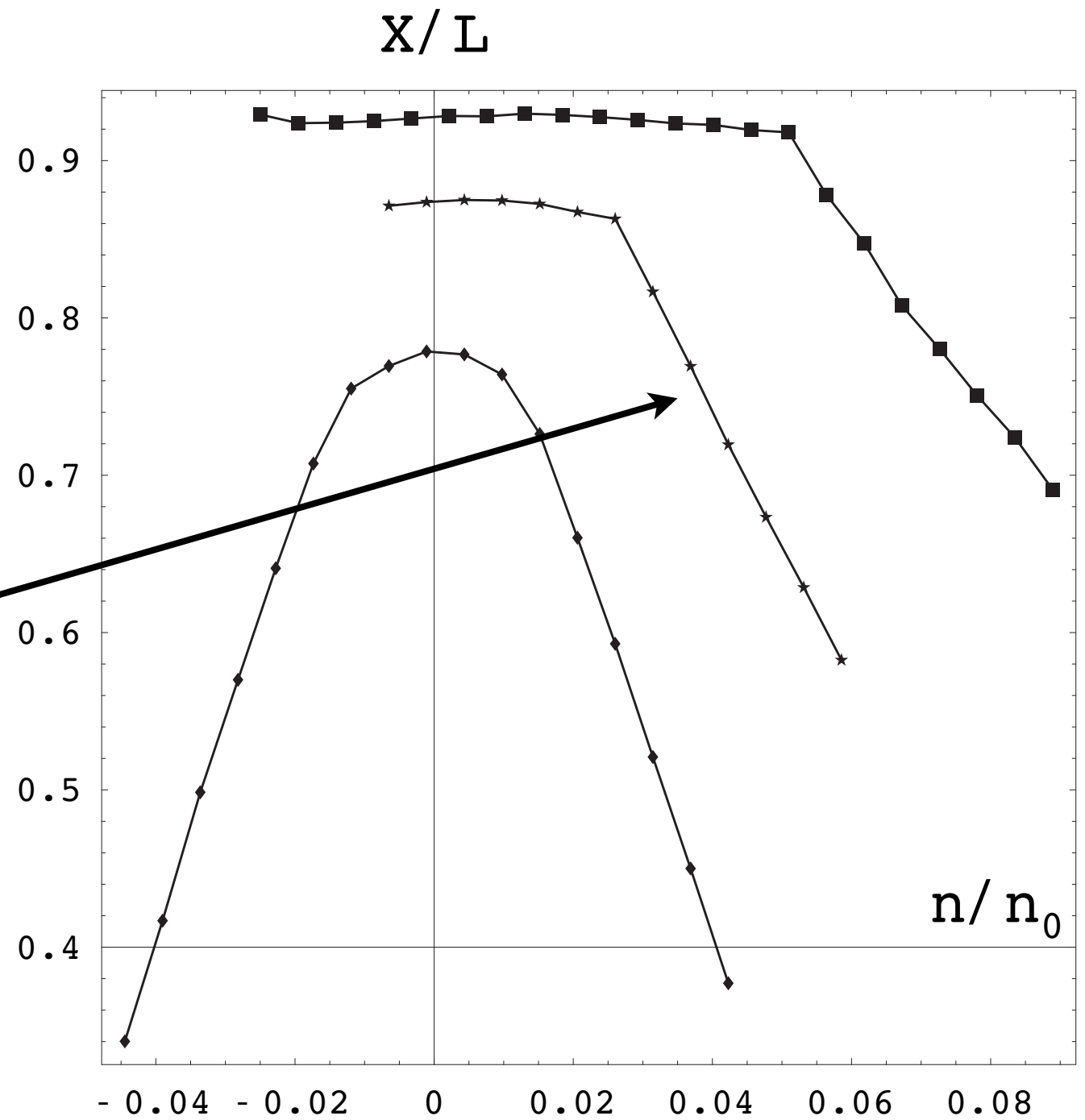
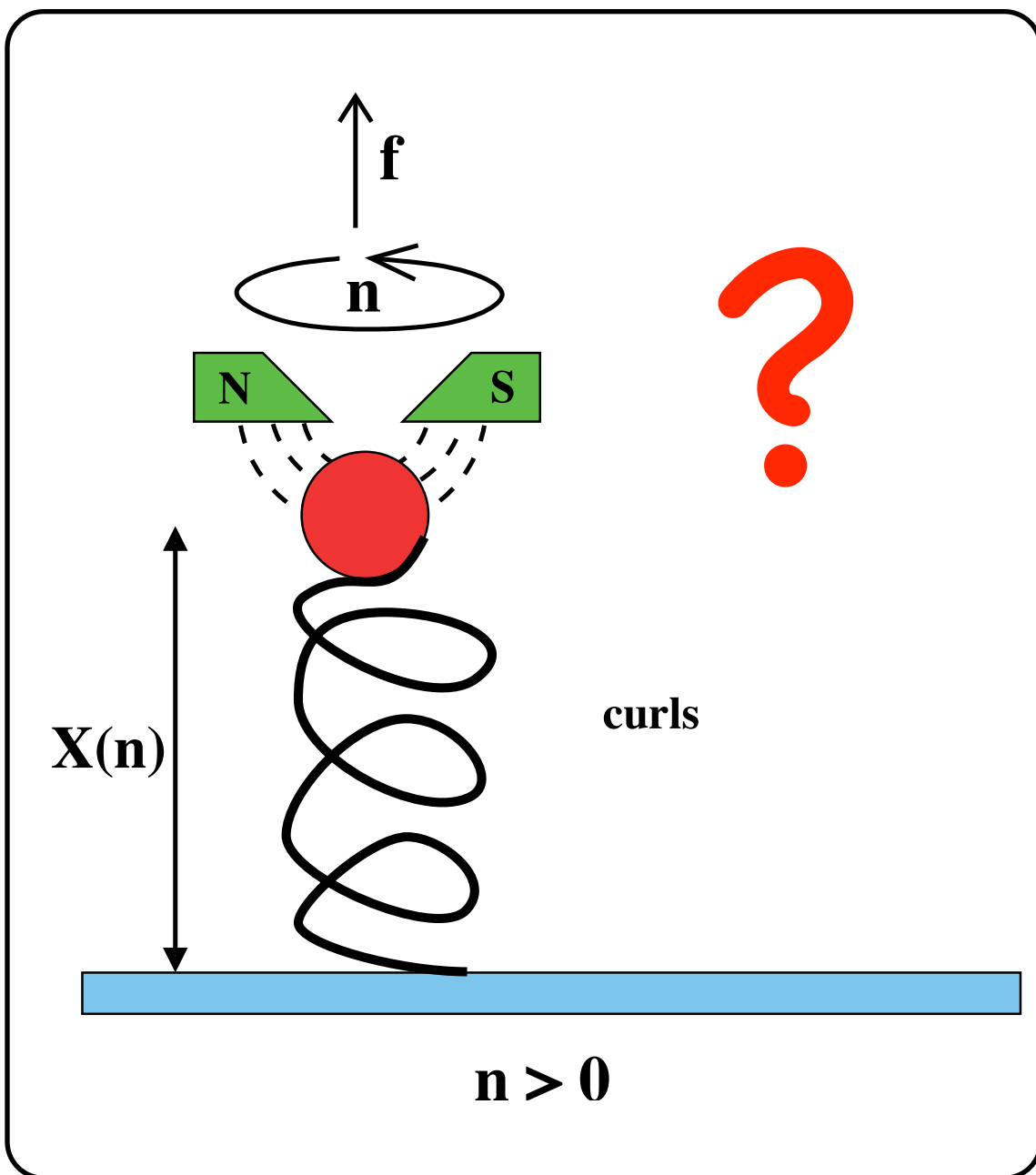
Data from G. Charvin (LPS-ENS)

Pulling and twisting DNA



Data from G. Charvin (LPS-ENS)

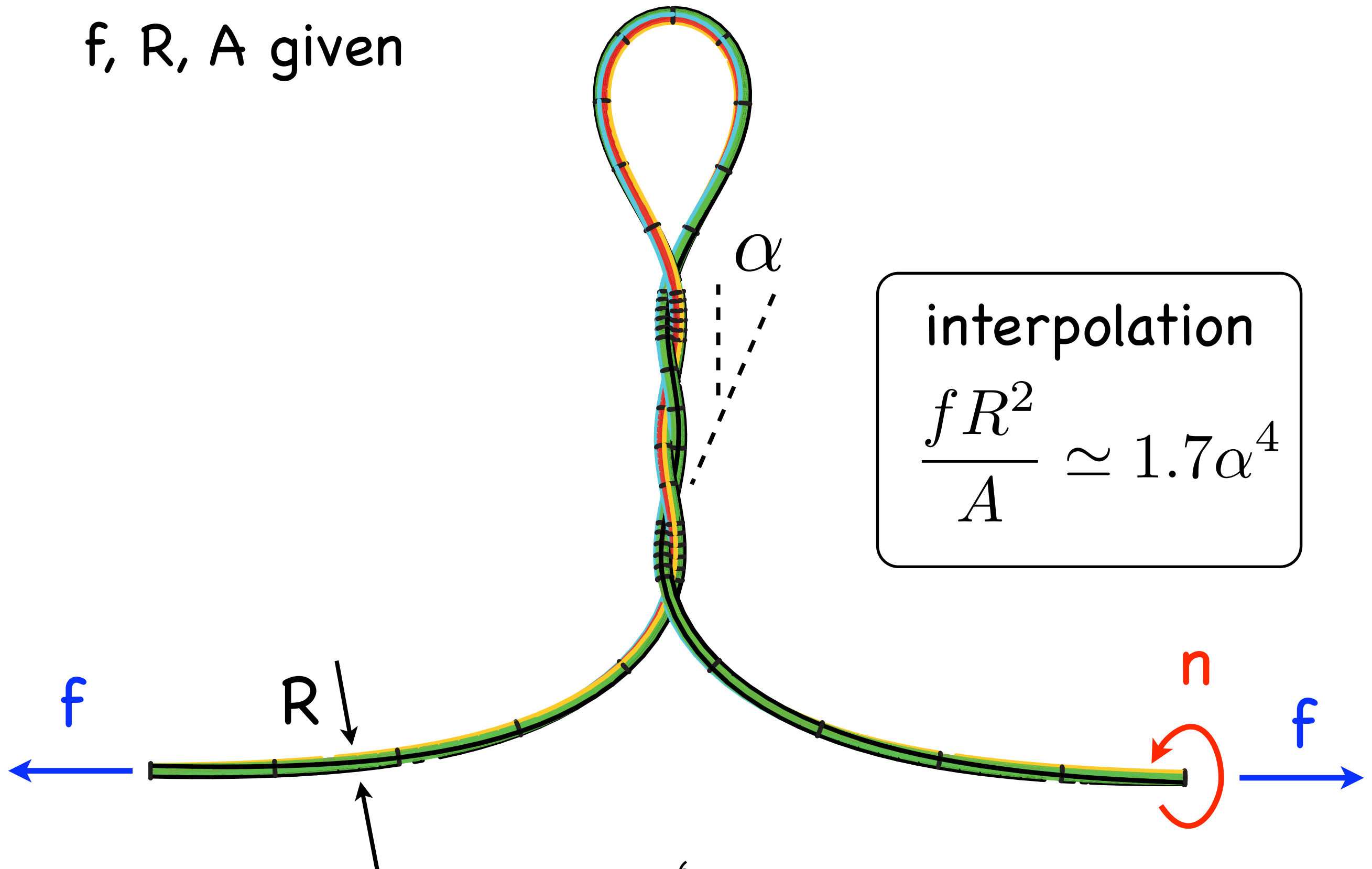
Pulling and twisting DNA



Data from G. Charvin (LPS-ENS)

Numerical simulations : twisted rod with self-contact

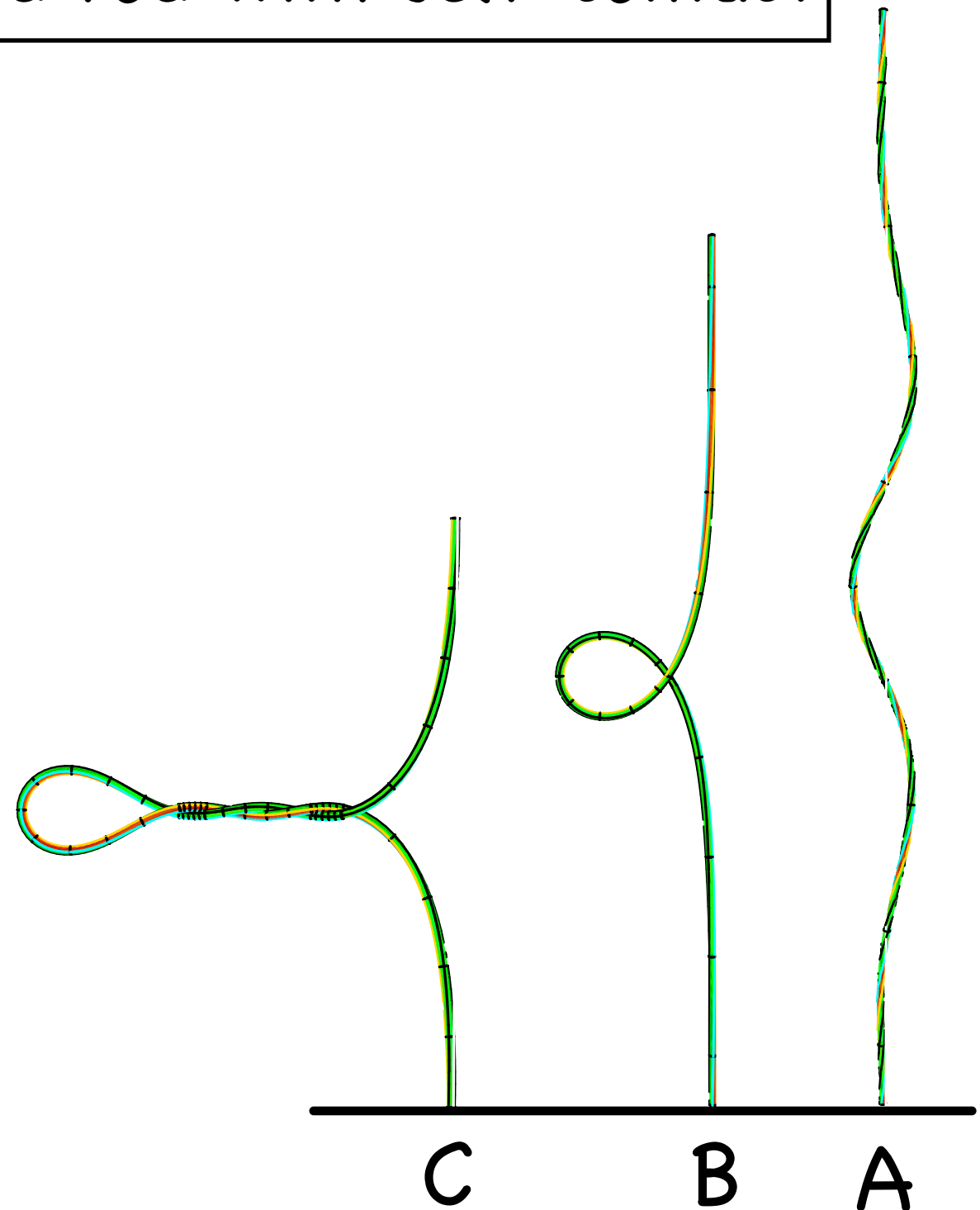
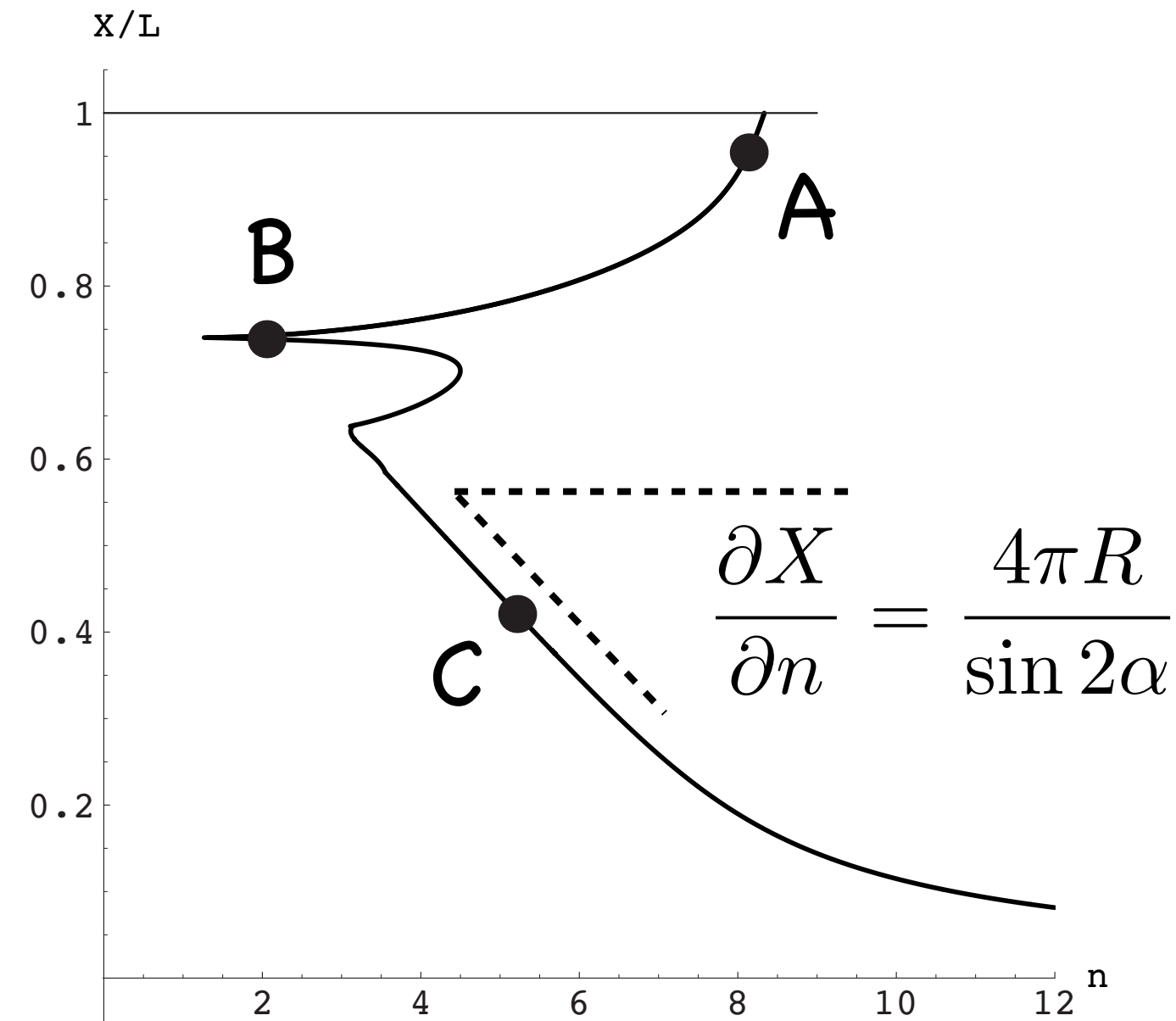
f, R, A given



interpolation

$$\frac{f R^2}{A} \simeq 1.7 \alpha^4$$

Numerical simulations : twisted rod with self-contact

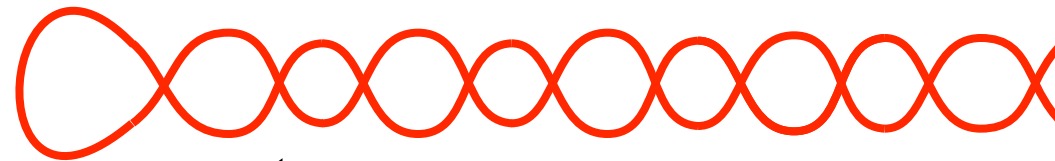


(based on Swigon+Coleman model for contact in Kirchhoff rods)

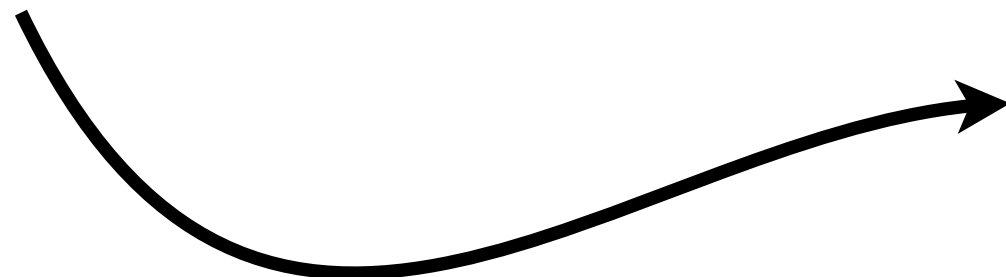
S. Neukirch, "Extracting DNA ... ", Phys. Rev. Lett. **93** (2004)

Limitations of the twisted rod model

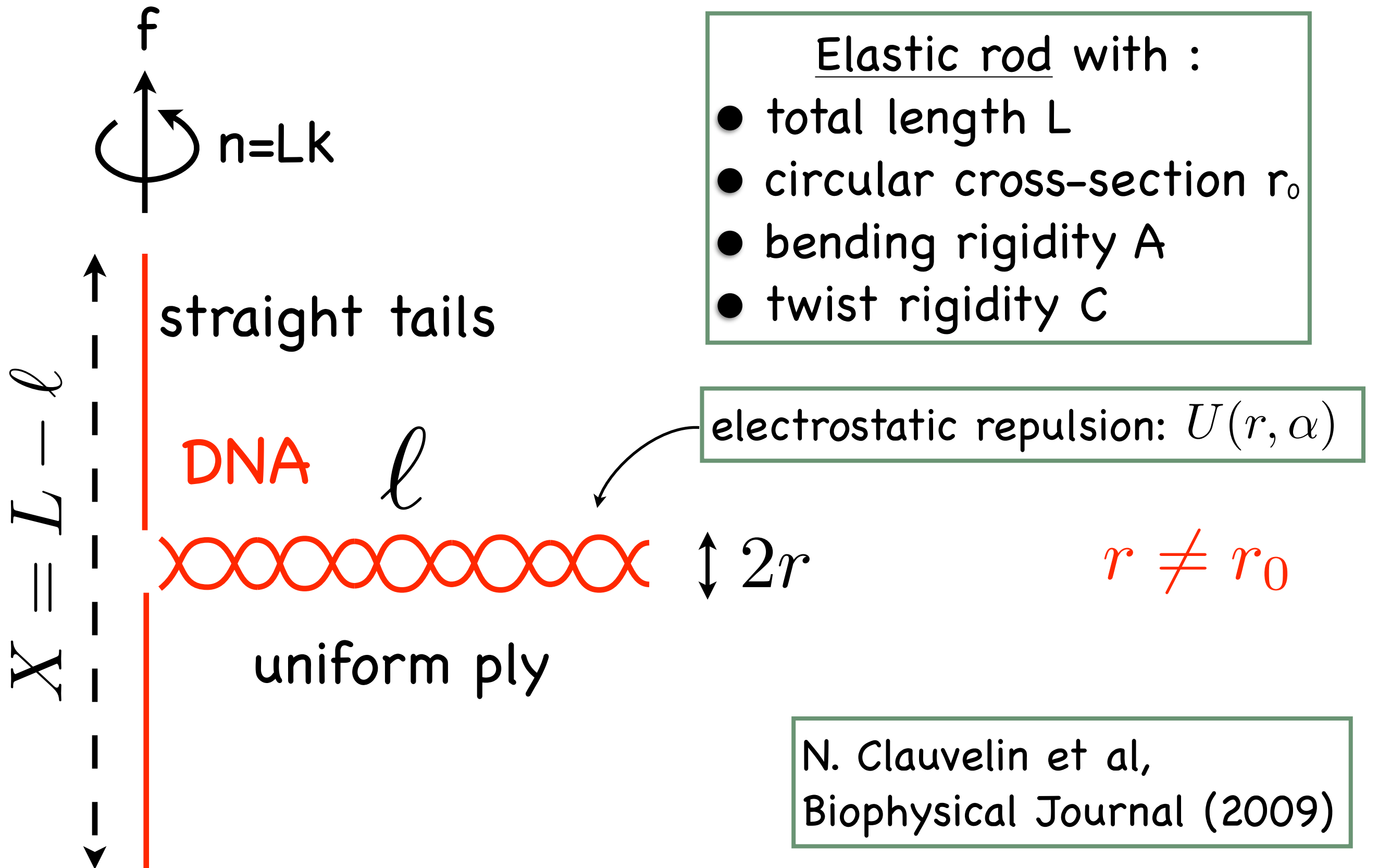
1. Electrostatics repulsion :
supercoiling radius R is not 1 nm



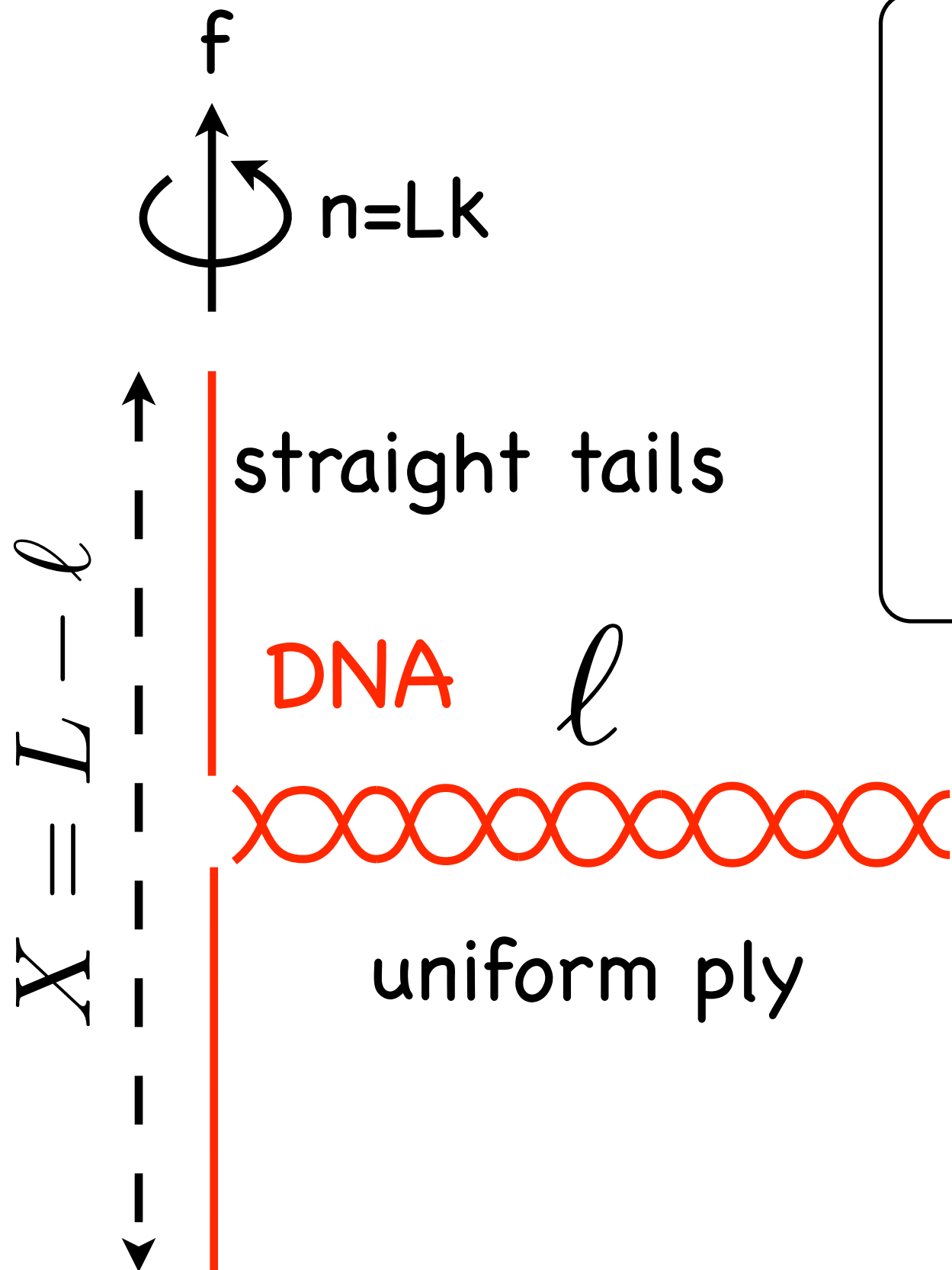
2. Tails are not straight :
«disordered walk» (Worm Like Chain)



Analytical model with electrostatics repulsion



Analytical model with electrostatics repulsion



$$\begin{aligned}
 V = & \frac{1}{2} A \frac{\sin^4 \alpha}{r^2} \ell && \text{bending} \\
 & + \frac{1}{2} C \tau^2 L && \text{twisting} \\
 & + U(R, \alpha) \ell && \text{electro} \\
 & - f X && \text{ext. force}
 \end{aligned}$$

with constraint:

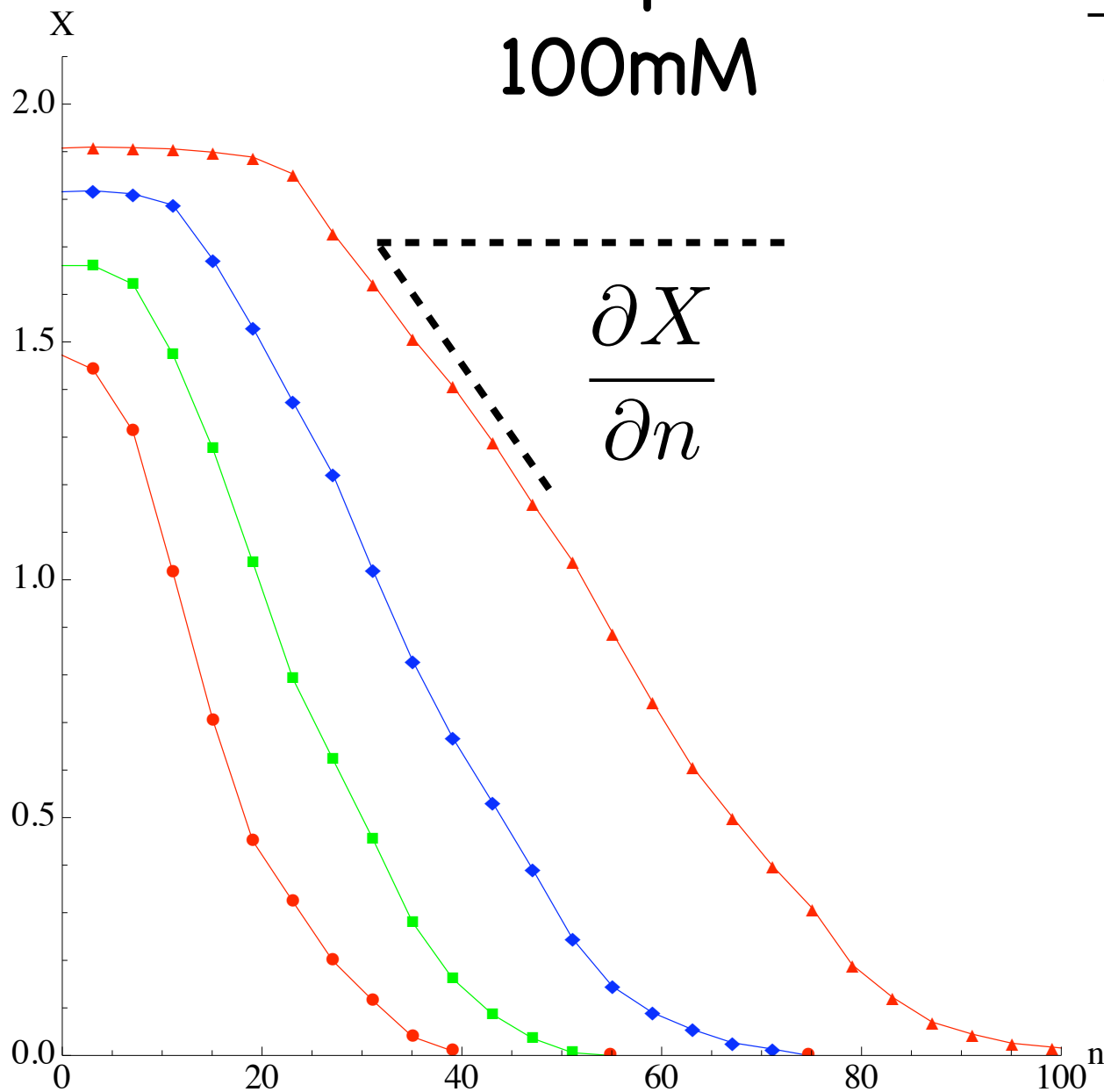
$$n = Lk = \frac{1}{2\pi} \left(\tau L + \frac{\sin 2\alpha}{2r} \ell \right)$$

minimize $V \Rightarrow (\alpha, r, \tau, \ell)$

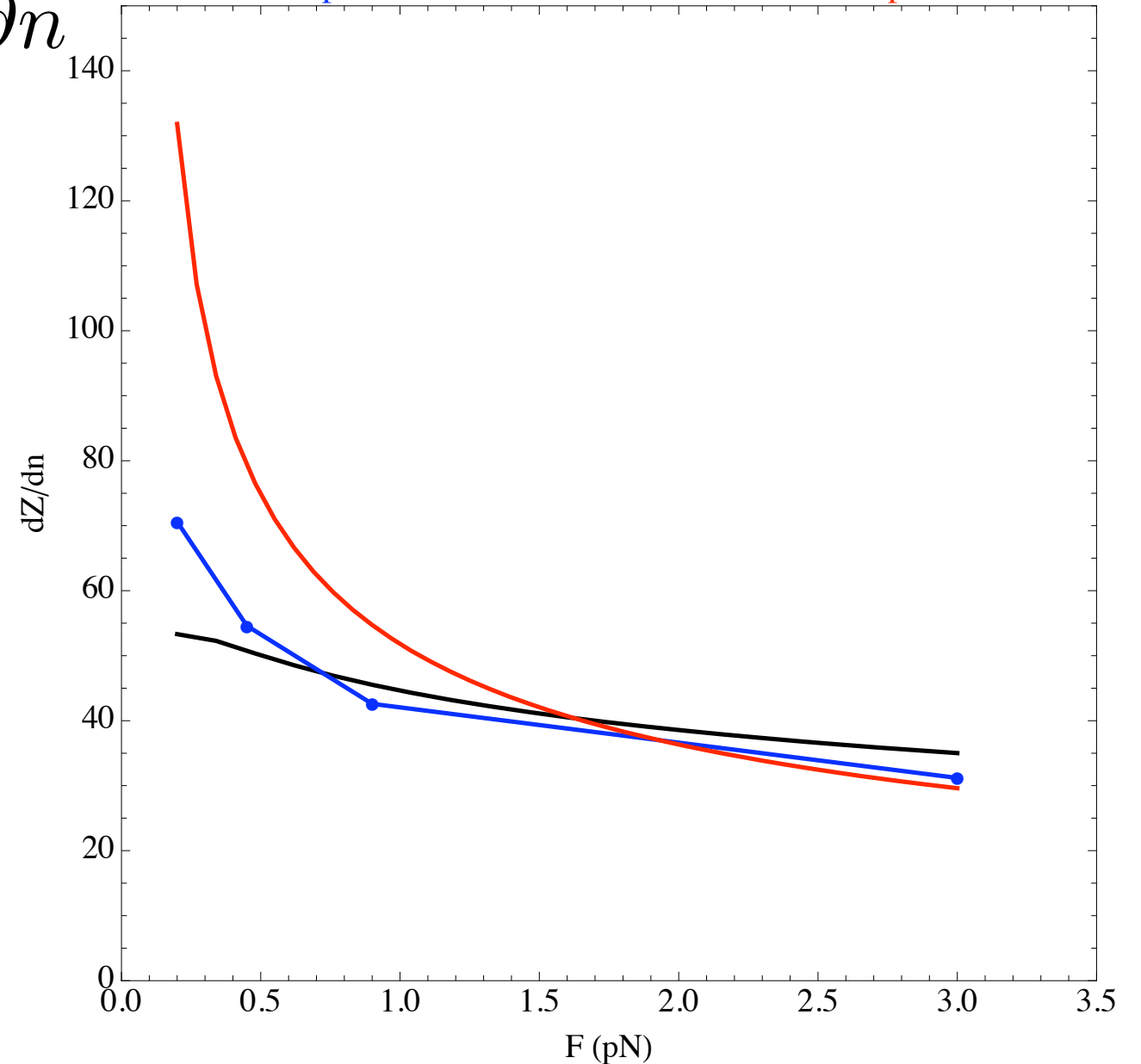
Analytical model with electrostatics repulsion : results

6kbp
100mM

$$\frac{\partial X}{\partial n}$$



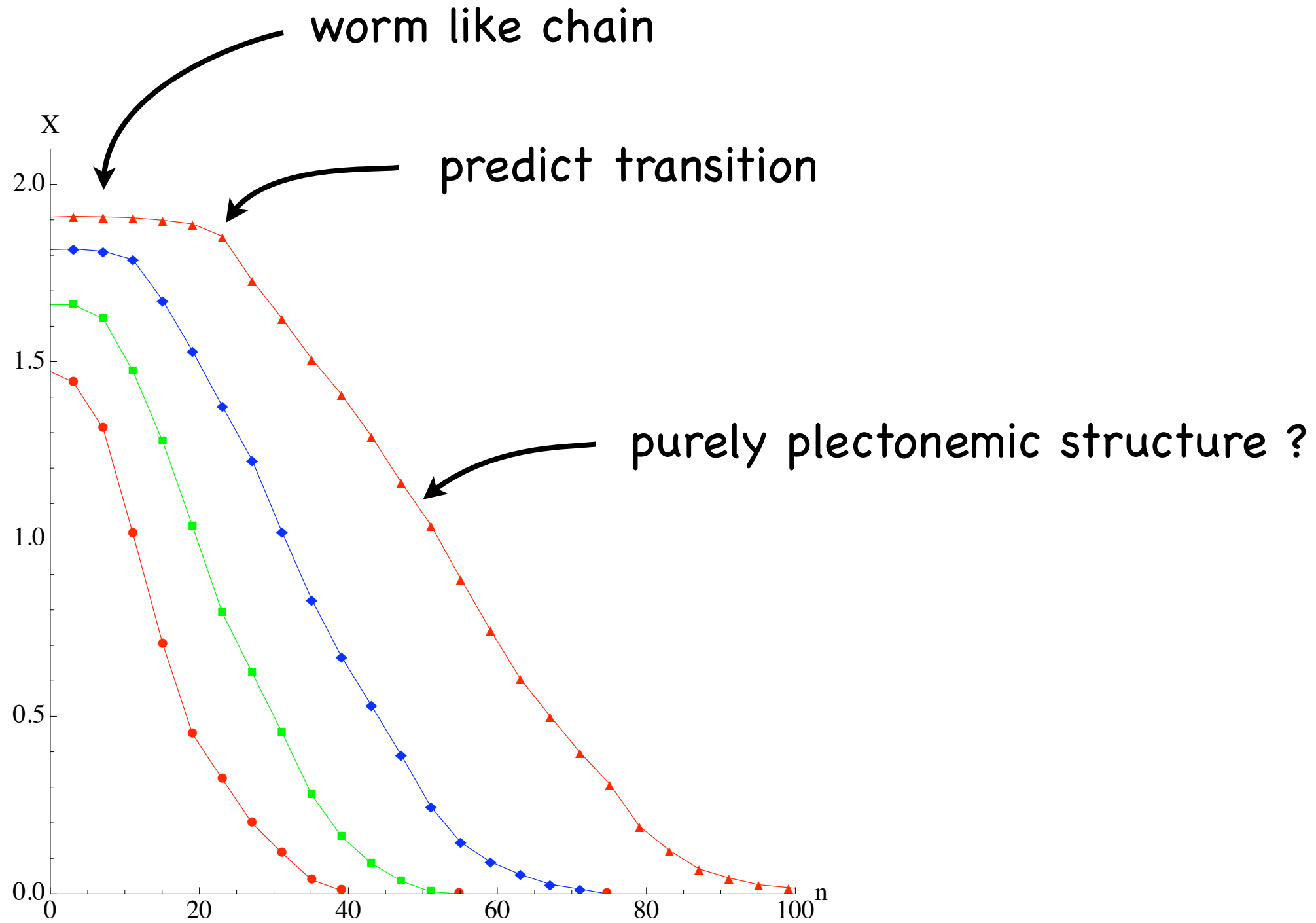
Experimental vs theoretical vs Marko slopes



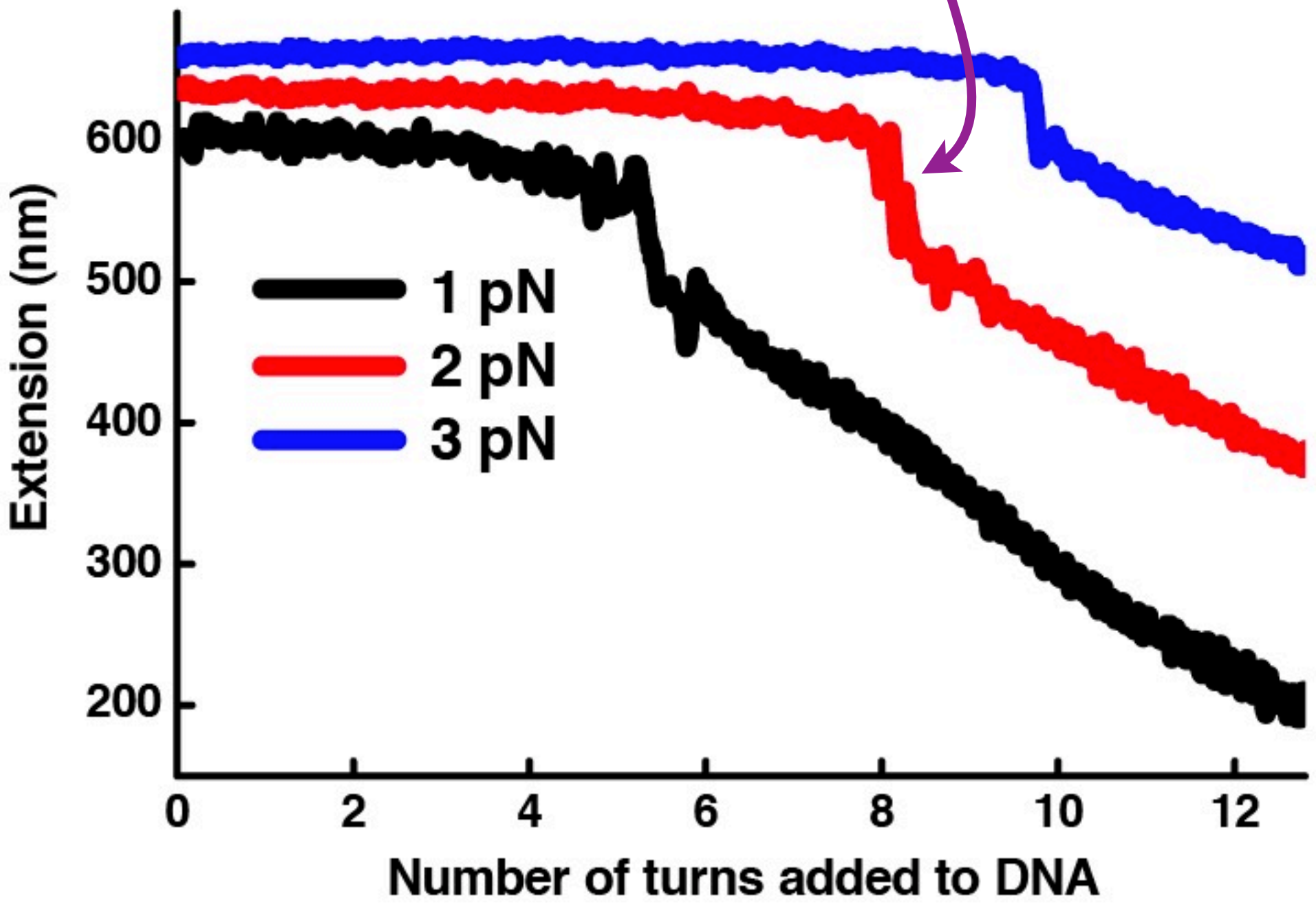
Data from G. Charvin (LPS-ENS)

J. Marko, "Torque and dynamics of linking number ...", Phys. Rev. E. (2007)

Statistical mechanics model



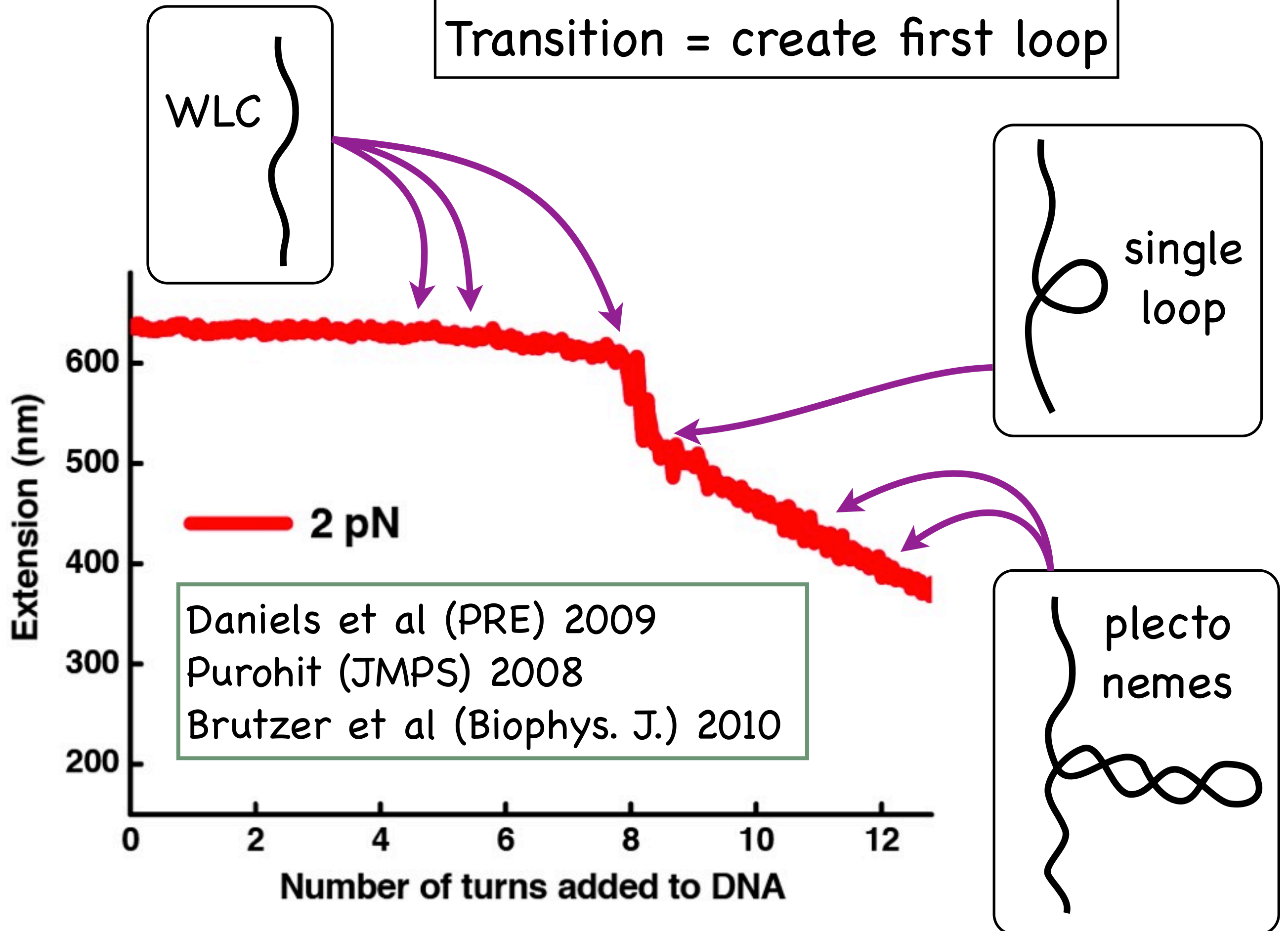
Abrupt transition



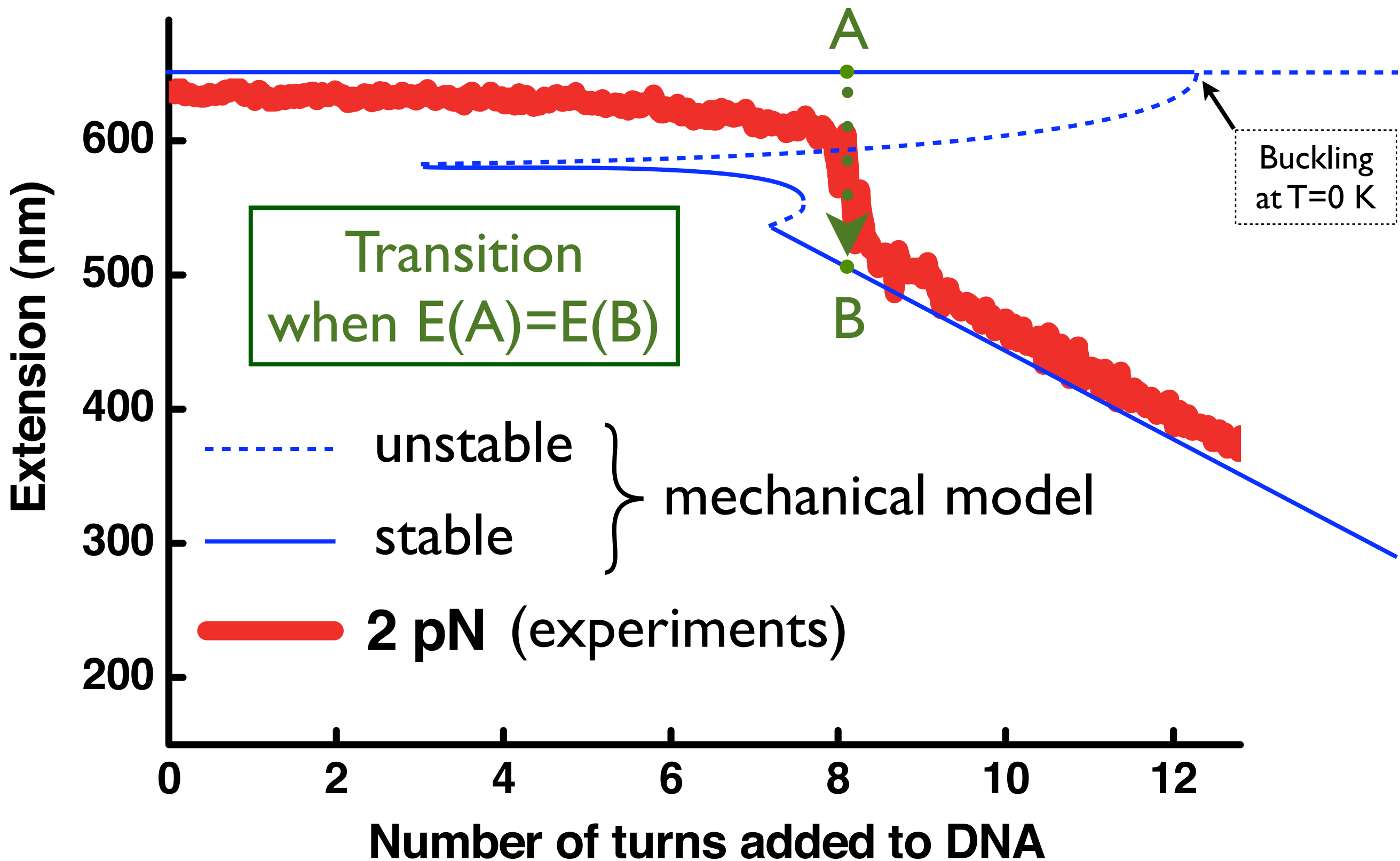
2.2 kbp
150 mM

Forth et al
Phys. Rev. Let.
2008

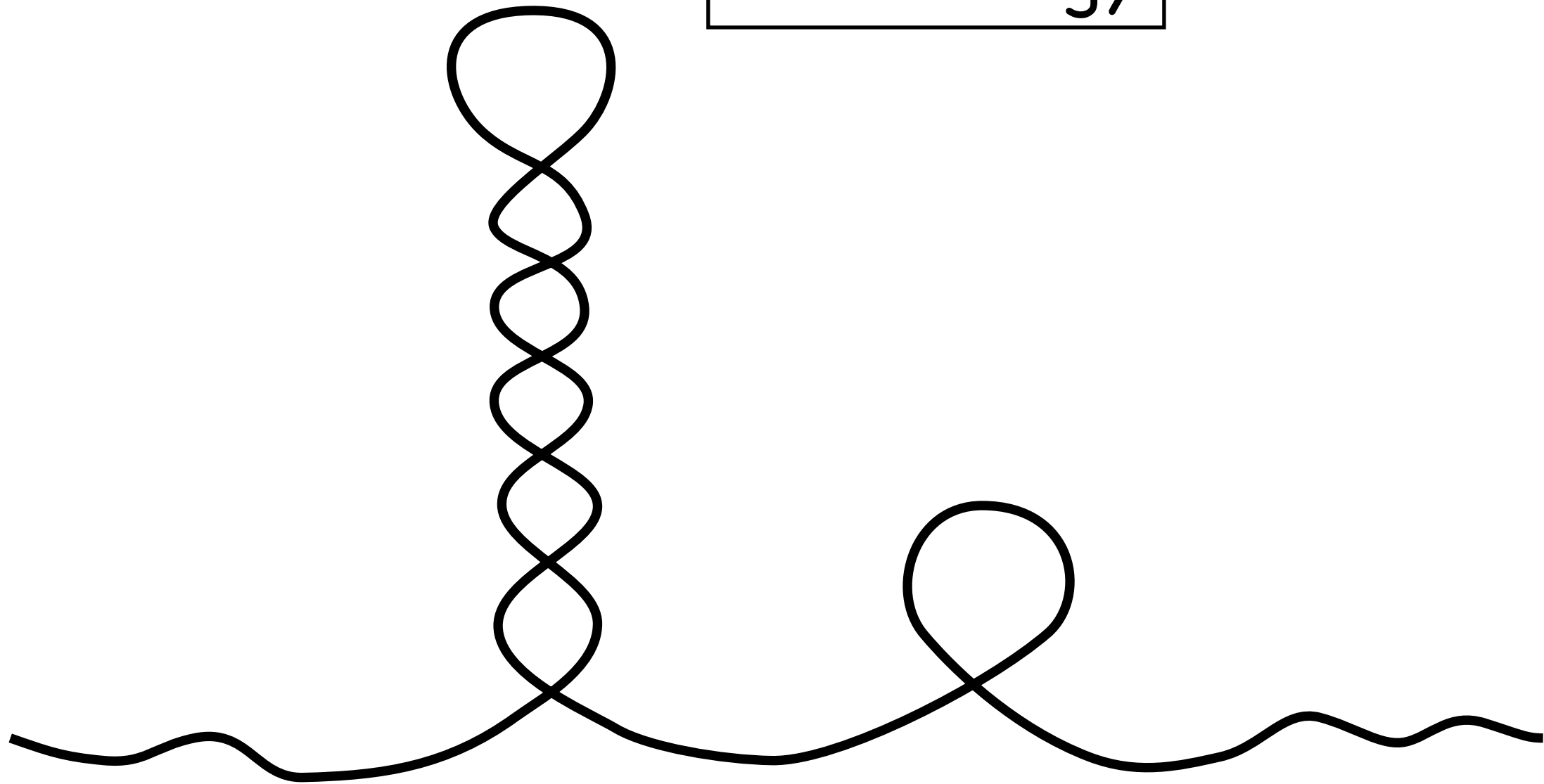
Transition = create first loop



Comparing free-energies of straight and supercoiled DNA

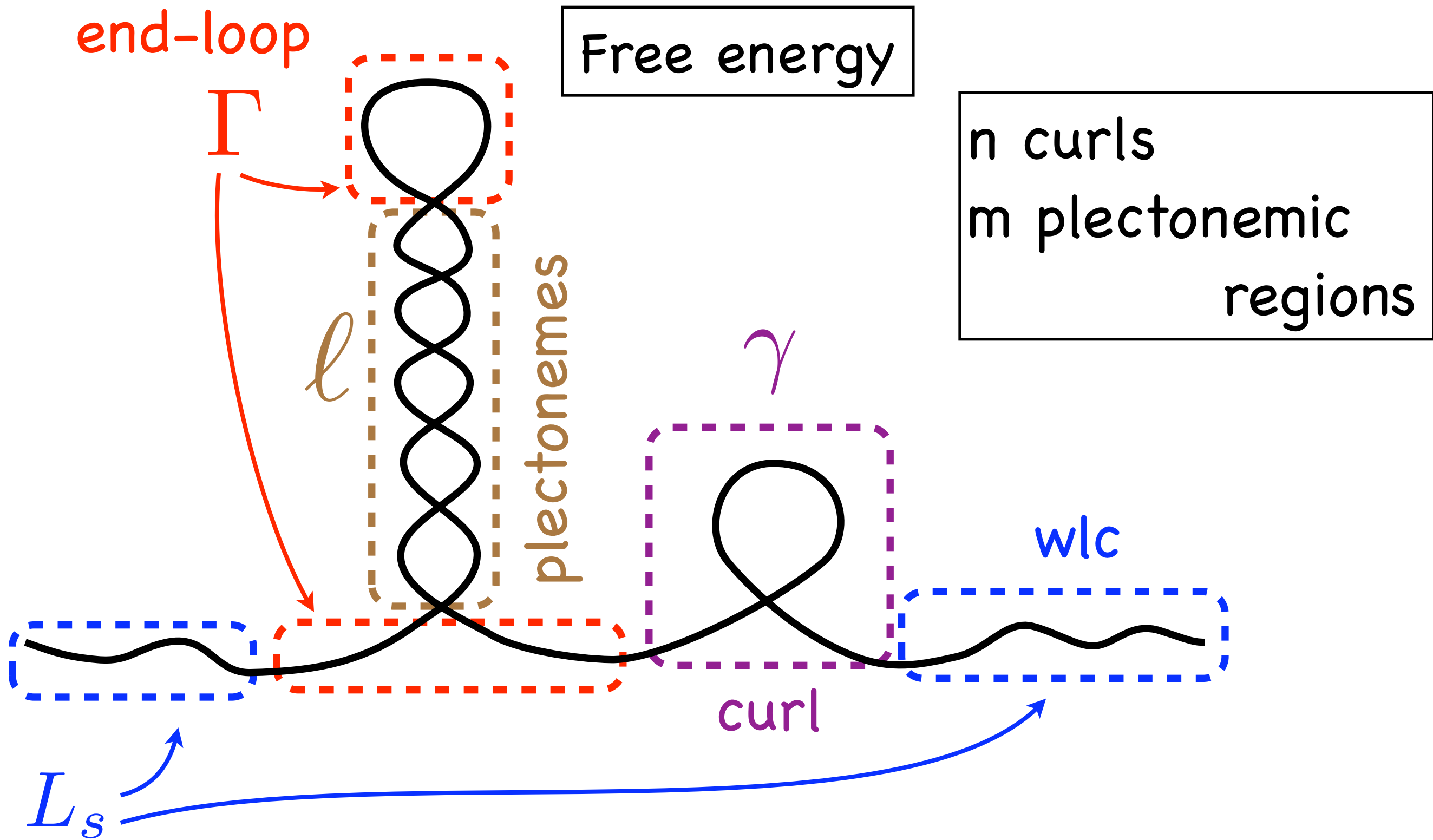


Free energy



total contour length

$$L = \text{nbp} \cdot 0.34 \text{ nm}$$



$$L = \text{nbp} \cdot 0.34 \text{ nm} = L_s + n\gamma + m(\Gamma + l)$$

Free energy

corrections to the
work of the
external force

$$L = L_s + n\gamma + m(\ell + \Gamma)$$

$$F = F_s + nF_\gamma + m(F_\Gamma + F_\ell) + F_f - TS$$

entropic terms

$$F_s = E_{twist} + E_{bend} + E_f = \frac{1}{2}C_s\tau_s^2 L_s - g(f)L_s$$

$$F_\gamma = E_{twist} + E_{bend} = \frac{1}{2}C\tau_p^2\gamma + 4\sqrt{Af}$$

$$F_\ell = E_{twist} + E_{bend} + E_{electro} = \frac{1}{2}C\tau_p^2\ell + \frac{1}{2}A\frac{\sin^4\alpha}{r^2}\ell + U\ell$$

$$F_\Gamma = E_{twist} + E_{bend} + E_{electro} = \frac{1}{2}C\tau_p^2\Gamma + q_b\sqrt{Af} + U\Gamma$$

$$F_f = 4n\sqrt{Af} + q_D m\sqrt{Af}$$

Free energy

$$L = L_s + n\gamma + m(\ell + \Gamma)$$

corrections to the work of the external force

$$F = F_s + nF_\gamma + m(F_\Gamma + F_\ell) + F_f - TS$$

entropic terms

$$U = U + \frac{1}{2} kT \left(\frac{kT}{Ar^2} \right)^{1/3} \quad \text{confinement in a tube}$$

$$F_\gamma = F_\gamma - \frac{1}{4} \text{Log} \left(\frac{4\pi^2 A (kT)^2}{d^4 f^3} \right) \quad \begin{array}{l} \text{loop size fluctuations} \\ (d=1\text{nm}) \end{array}$$

$$S(L, \ell, n, m) = (n + m) \text{Log} \left(\frac{L - m\ell}{\gamma} - n - m \right) - \text{Log} n! - \text{Log} m!$$

Tonks hard core gas

Free energy

$$F = F_s + n F_\gamma + m (F_\Gamma + F_\ell) + F_f - T S$$

study F under the two constraints:

$$L = L_s + n\gamma + m(\ell + \Gamma)$$

$$n = Lk = \underbrace{\frac{\tau_s}{2\pi} L_s + \frac{\tau_p}{2\pi} (m\ell + n\gamma + m\Gamma)}_{\text{Twist}} + \underbrace{\frac{\sin 2\alpha}{4\pi r} m\ell + mp + n\Lambda}_{\text{Writhe}}$$

we replace τ_s, L_s to obtain :

$$F_{nm} = F_{nm}(\alpha, r, \tau_p, \ell)$$

Computation method

$$m \neq 0 \quad F_{nm} = F_{nm}(\alpha, r, \tau_p, \ell)$$

$$m = 0 \quad F_{n0} = F_{n0}(\alpha, r, \tau_p)$$

$$\text{reference state} \quad F_{00} = -gL + \frac{1}{2}C_s \left(\frac{2\pi Lk}{L} \right)^2 L$$

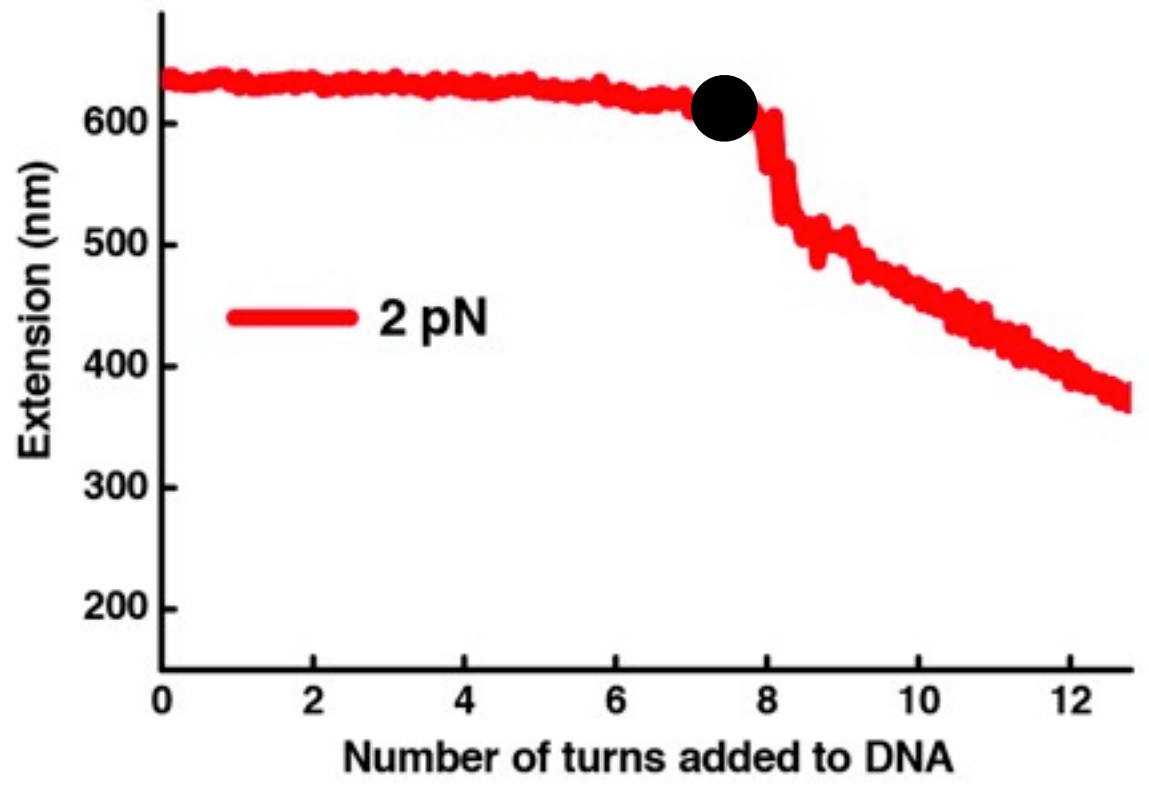
$$\Delta F_{nm}(\alpha, r, \tau_p, \ell) = F_{nm} - F_{00}$$

minimize with regard to α, r, τ_p (saddle point approx.)

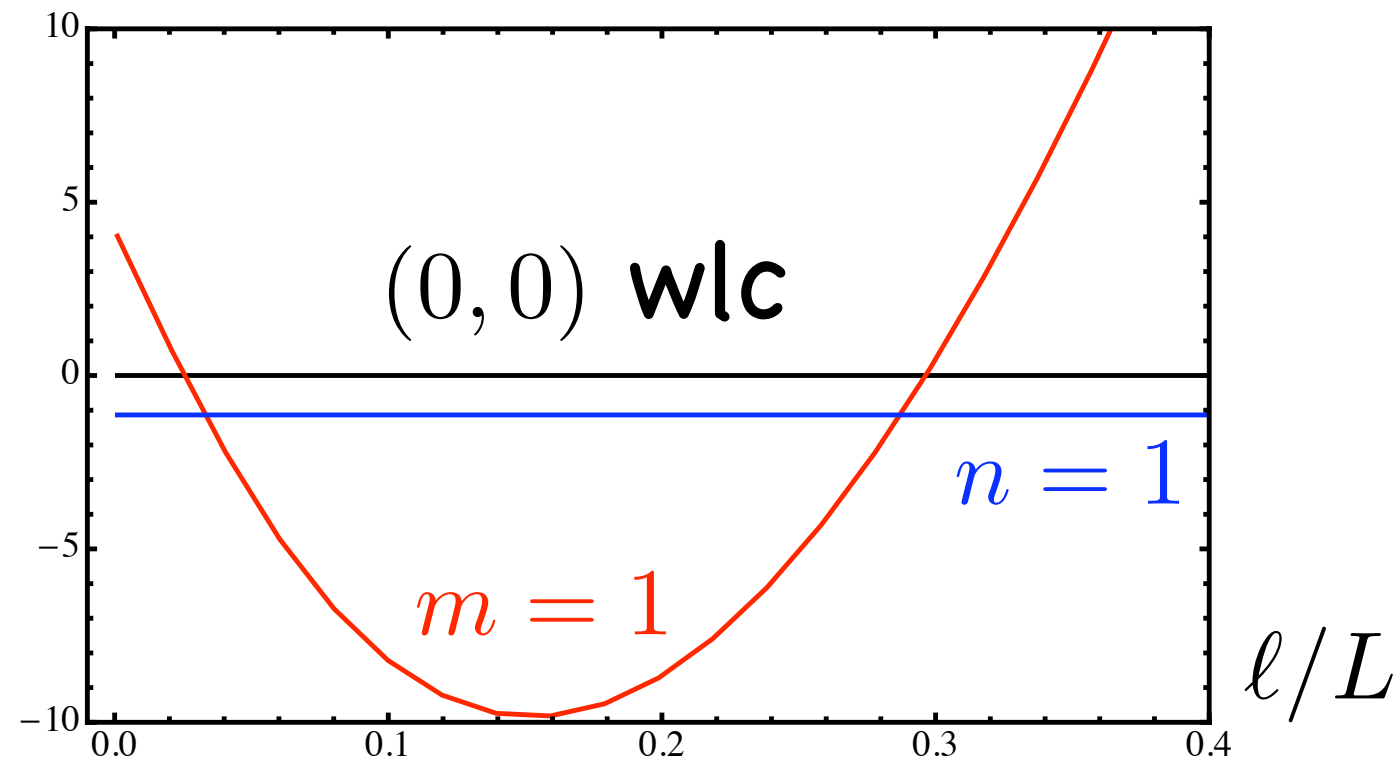
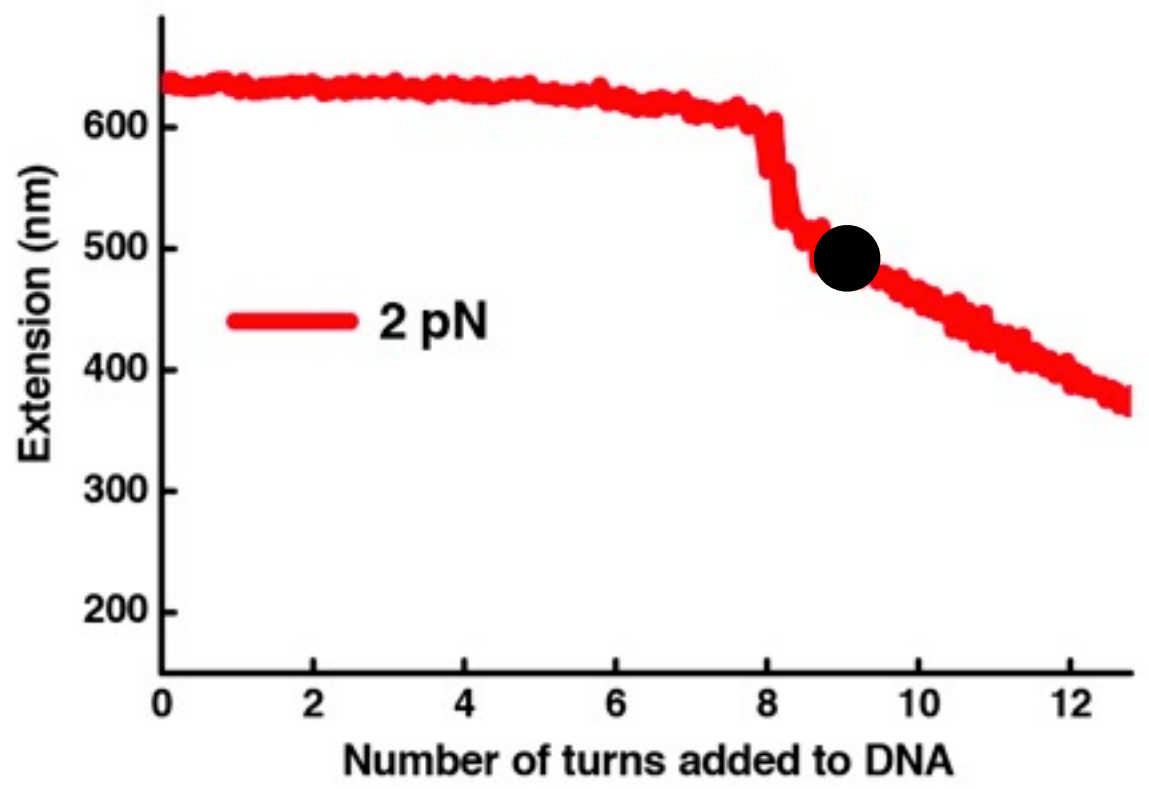
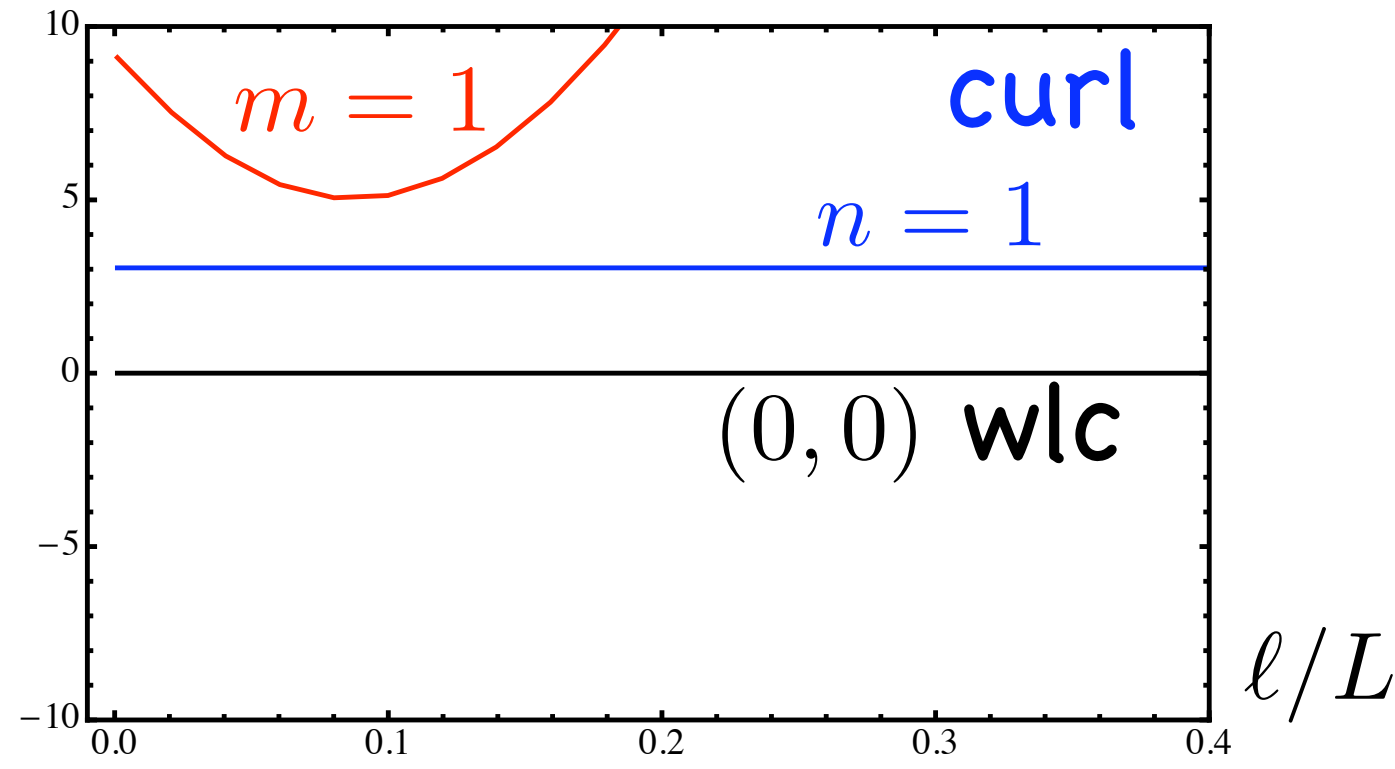
plot as function of ℓ

Results

Experimental data



$\Delta F/kT$ plecto



Computing mean values

At fixed Lk

Probability of seeing a configuration with

n curls and

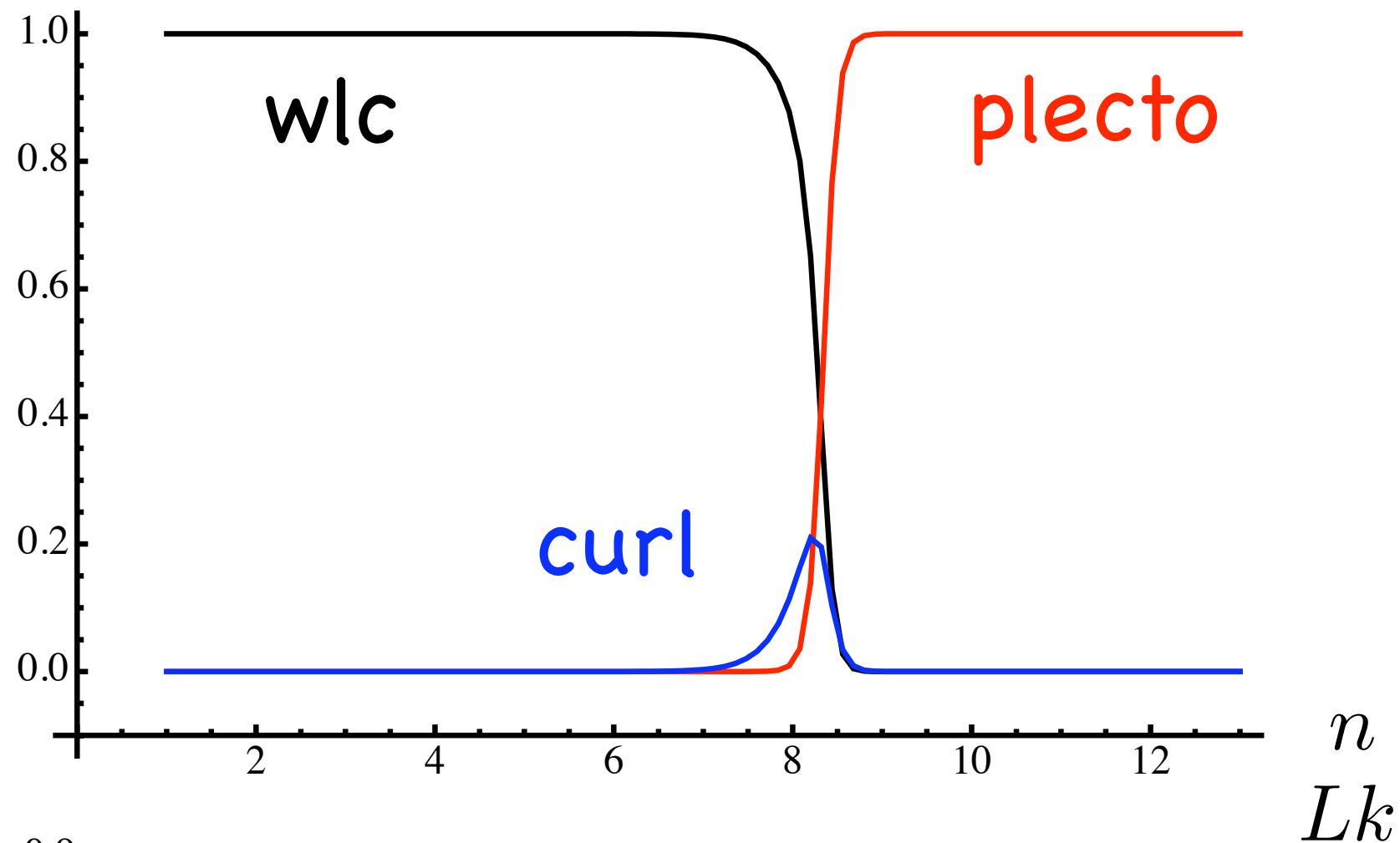
m plectonemic parts: $P_{nm}(n = Lk) = \frac{\int e^{-F_{nm}(\ell)} d\ell}{\sum_{n,m} \int e^{-F_{nm}(\ell)} d\ell}$

Mean extension of the system (considering all possible configurations:

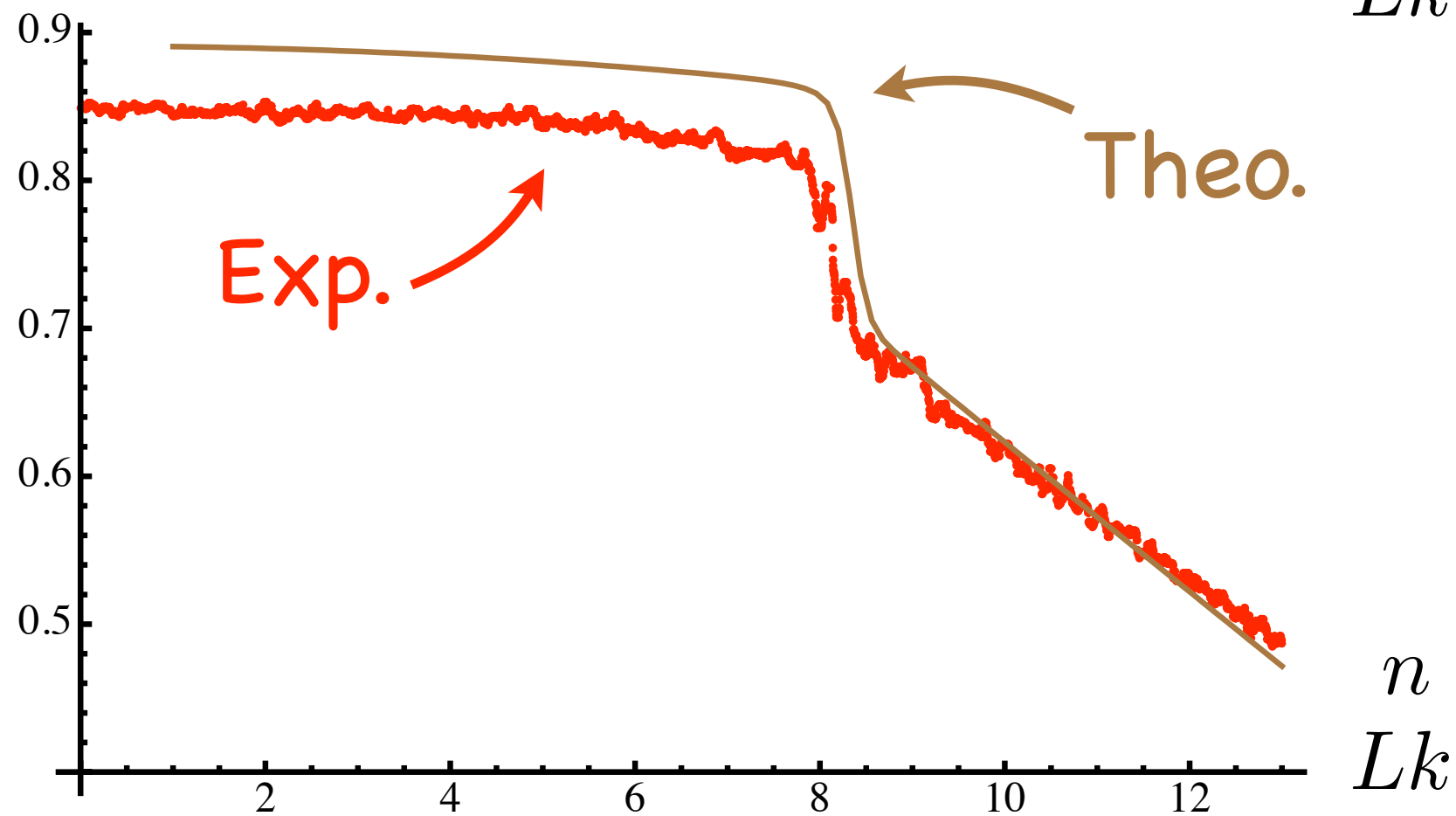
$$\langle X \rangle (n = Lk) = \frac{\sum_{n,m} \int X_{nm}(\ell) e^{-F_{nm}(\ell)} d\ell}{\sum_{n,m} \int e^{-F_{nm}(\ell)} d\ell}$$

Results

P_{nm}



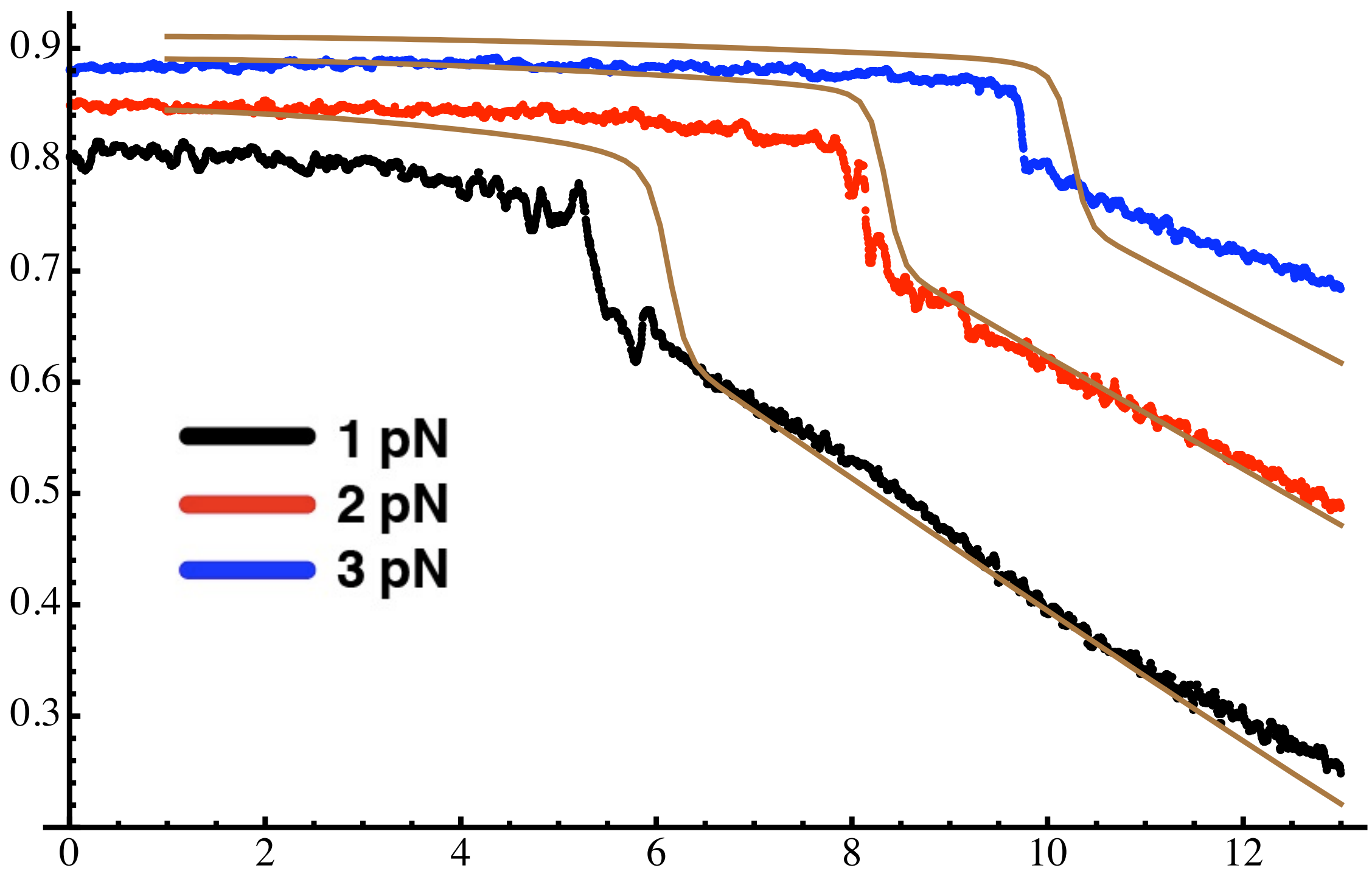
$\langle X/L \rangle$



2.2 kbp
150 mM
2 pN

data : Forth et al
(Phys. Rev. Let.) 2008

Results

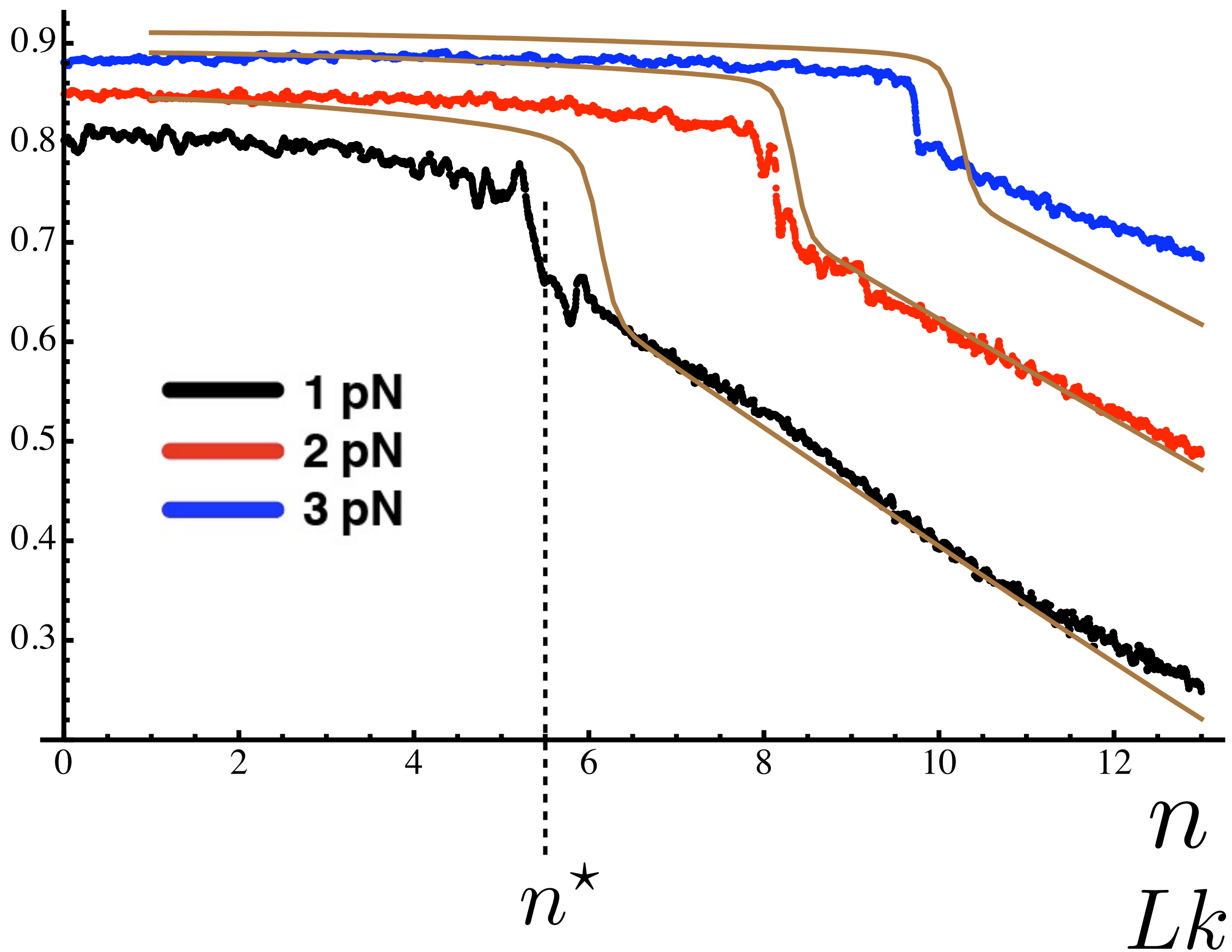


2.2 kbp
150 mM

data : Forth et al
(Phys. Rev. Let.) 2008

n
 Lk

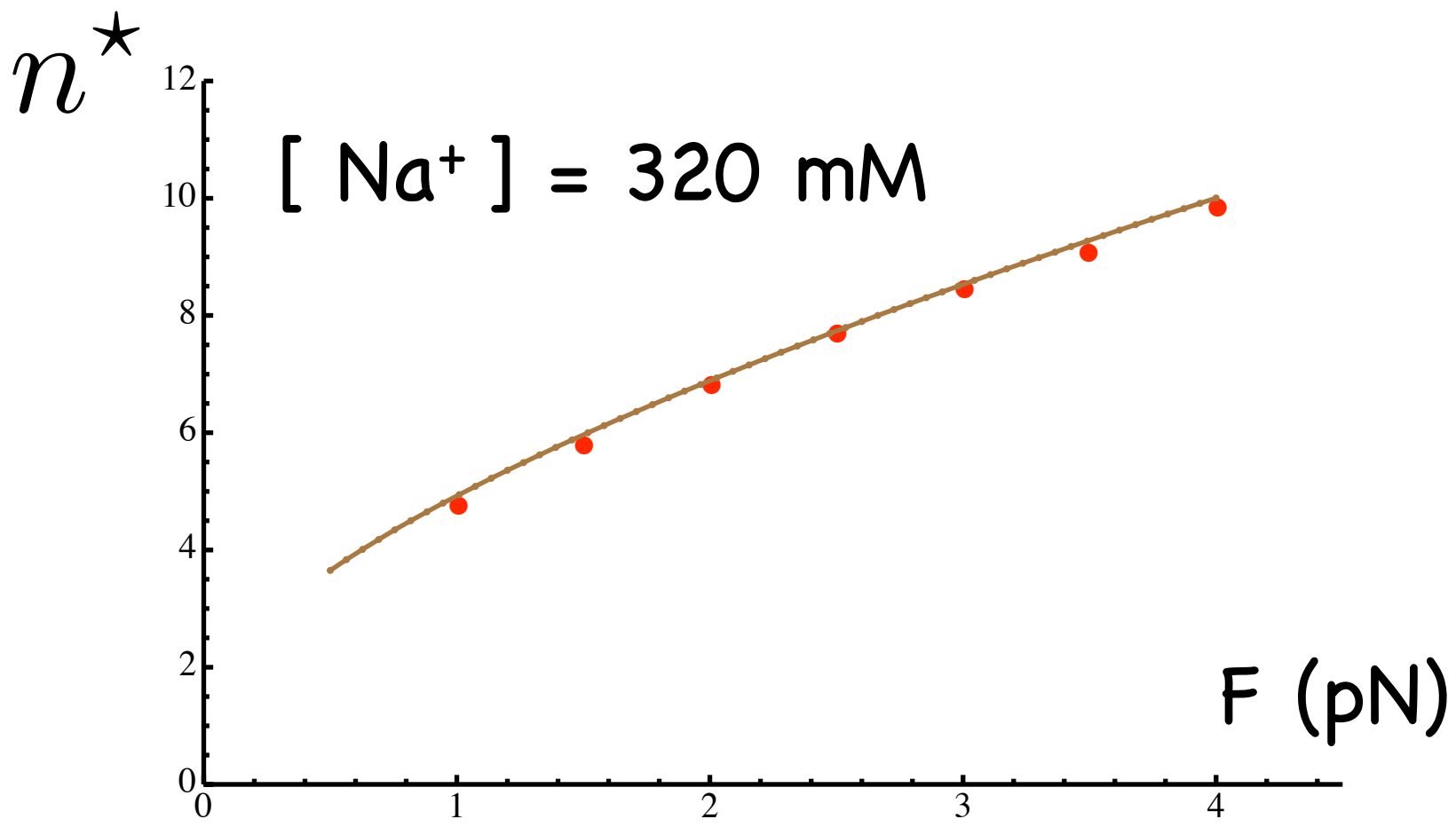
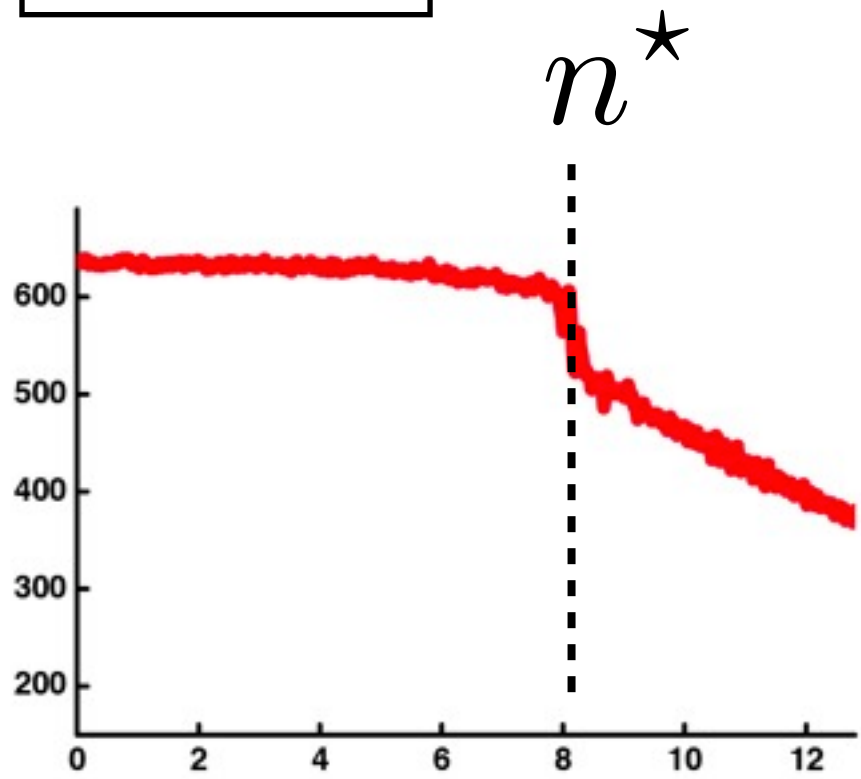
Results



2.2 kbp
150 mM

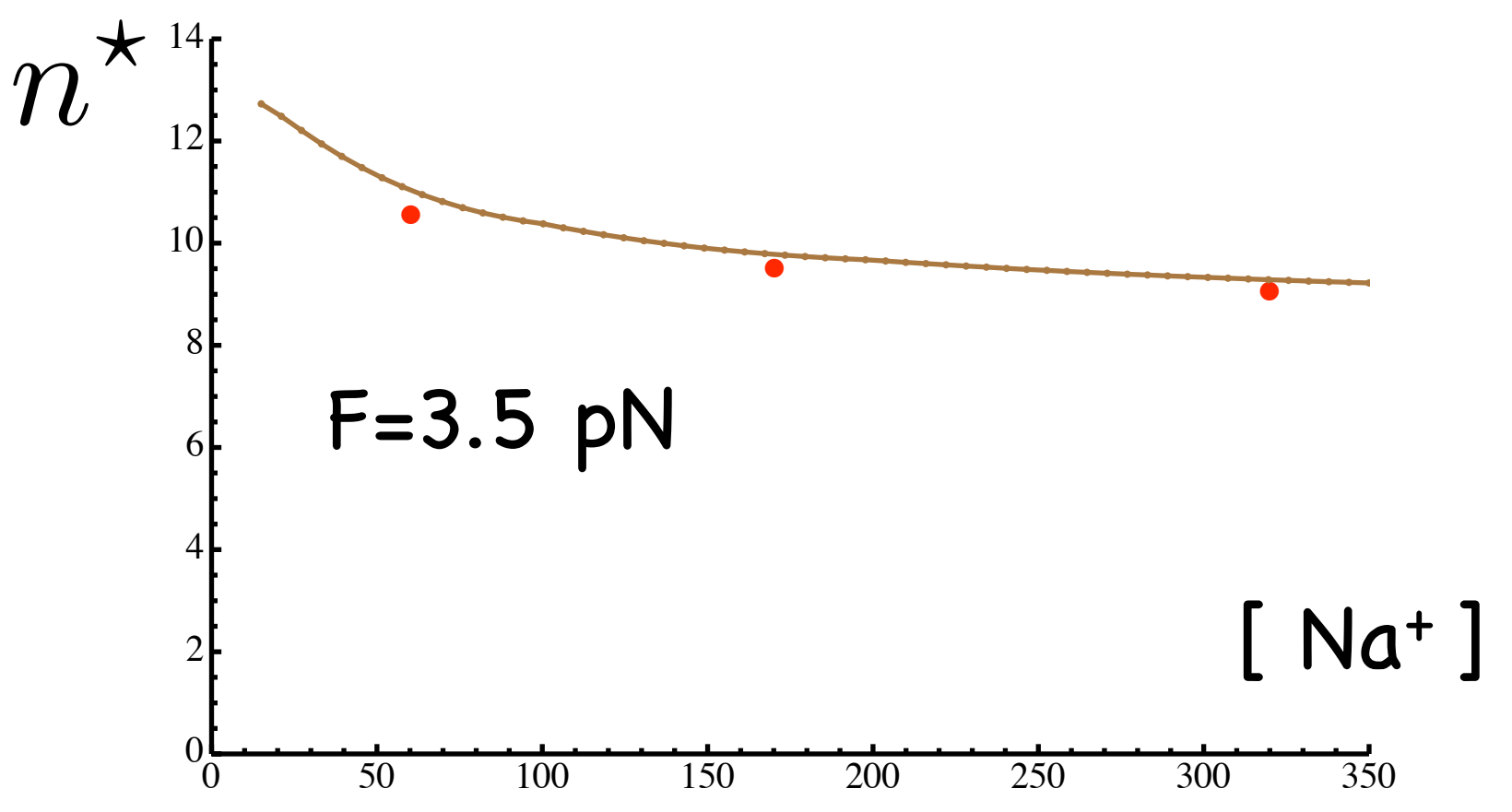
data : Forth et al
(Phys. Rev. Let.) 2008

Results



1900 bp

data : Brutzer et al
(Biophys. J.) 2010



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