

Elasticity and Electrostatics of plectonemic DNA

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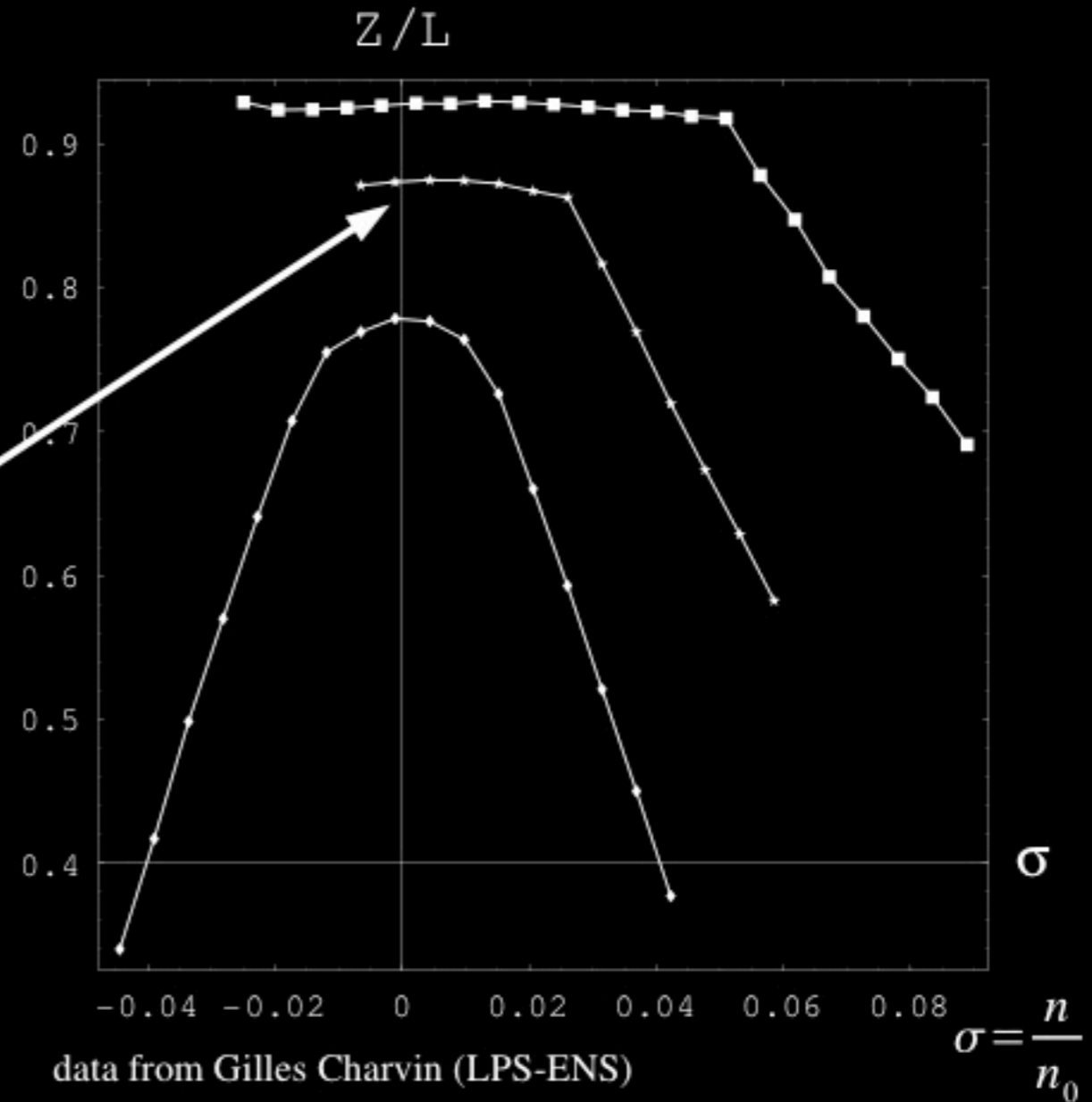
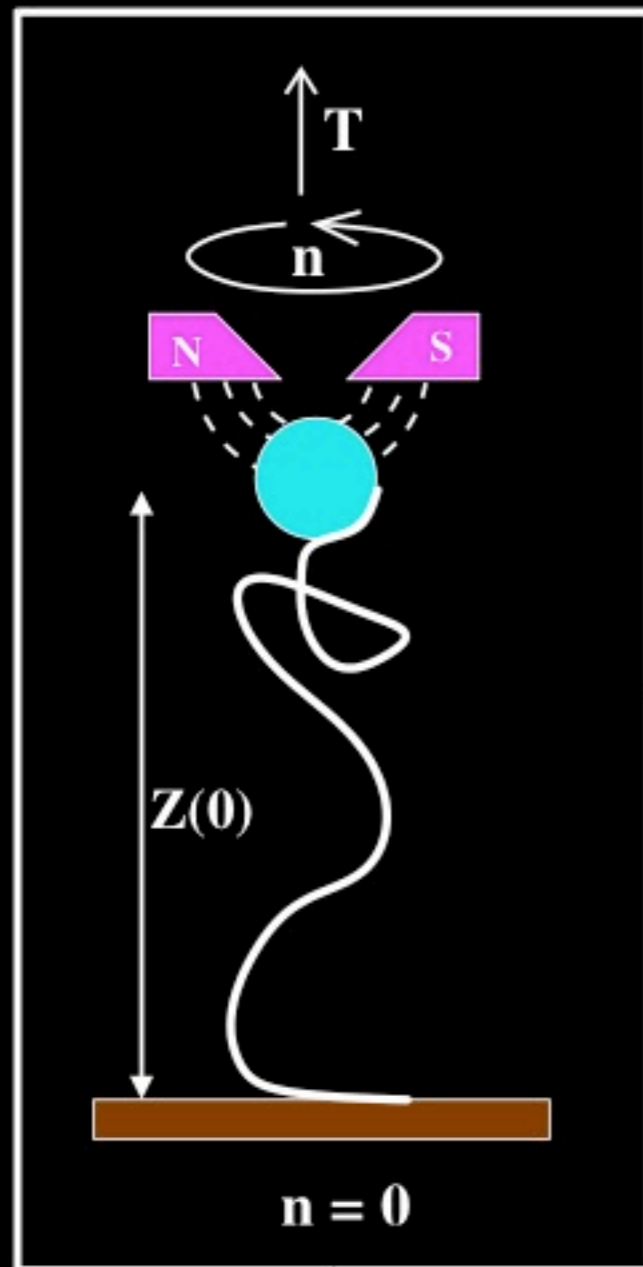
joint work with:
Nicolas Clauvelin (PhD work)
Basile Audoly

Why study DNA mechanical properties ?

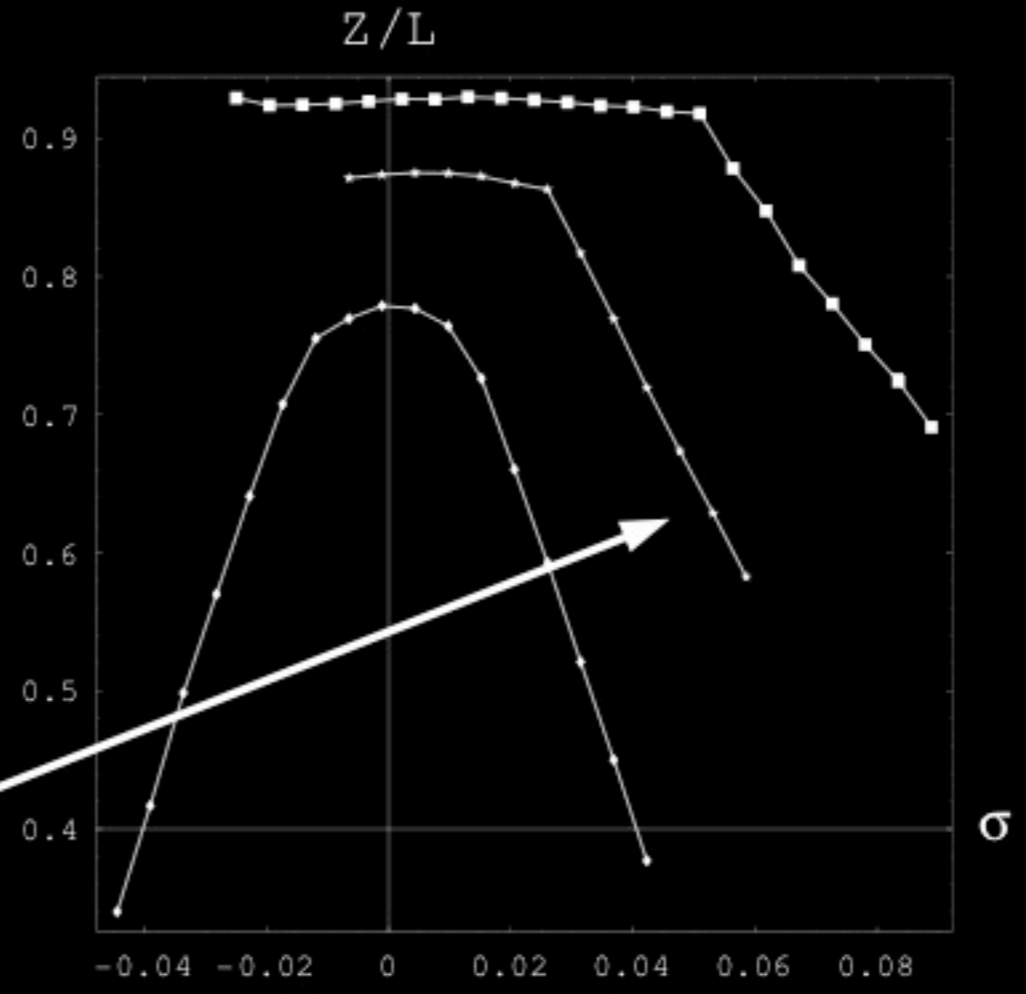
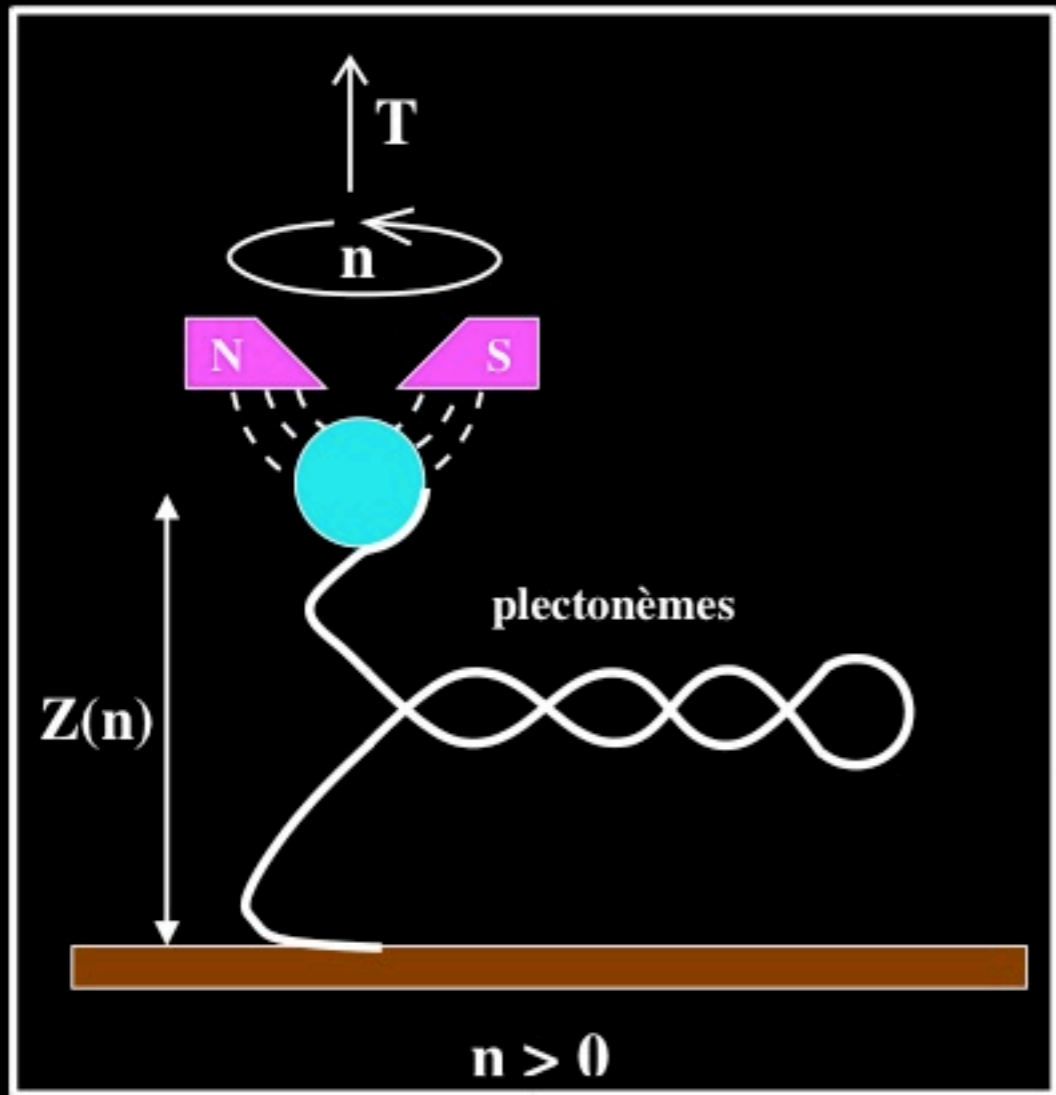
mechanical properties influence biology of the cell

- 2 meters of DNA in a 10 micron wide nucleus
- ejection from viral capsid
- transcription (RNAPolymerase is torque dependent)
- protein binding is strain dependent, or induces strain on DNA
- chromatin compaction/decompaction (cell division)

Pulling and twisting DNA



Pulling and twisting DNA

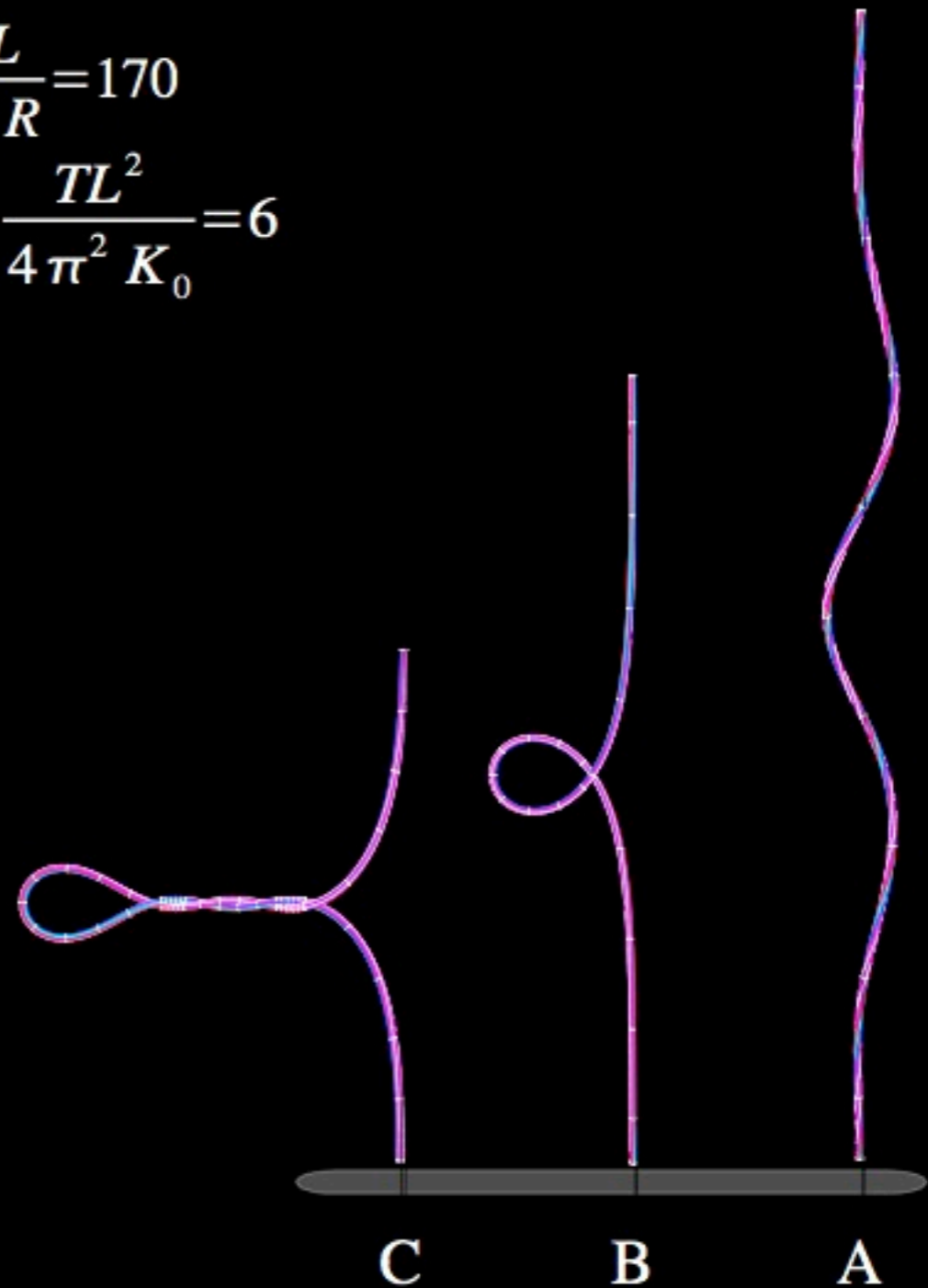
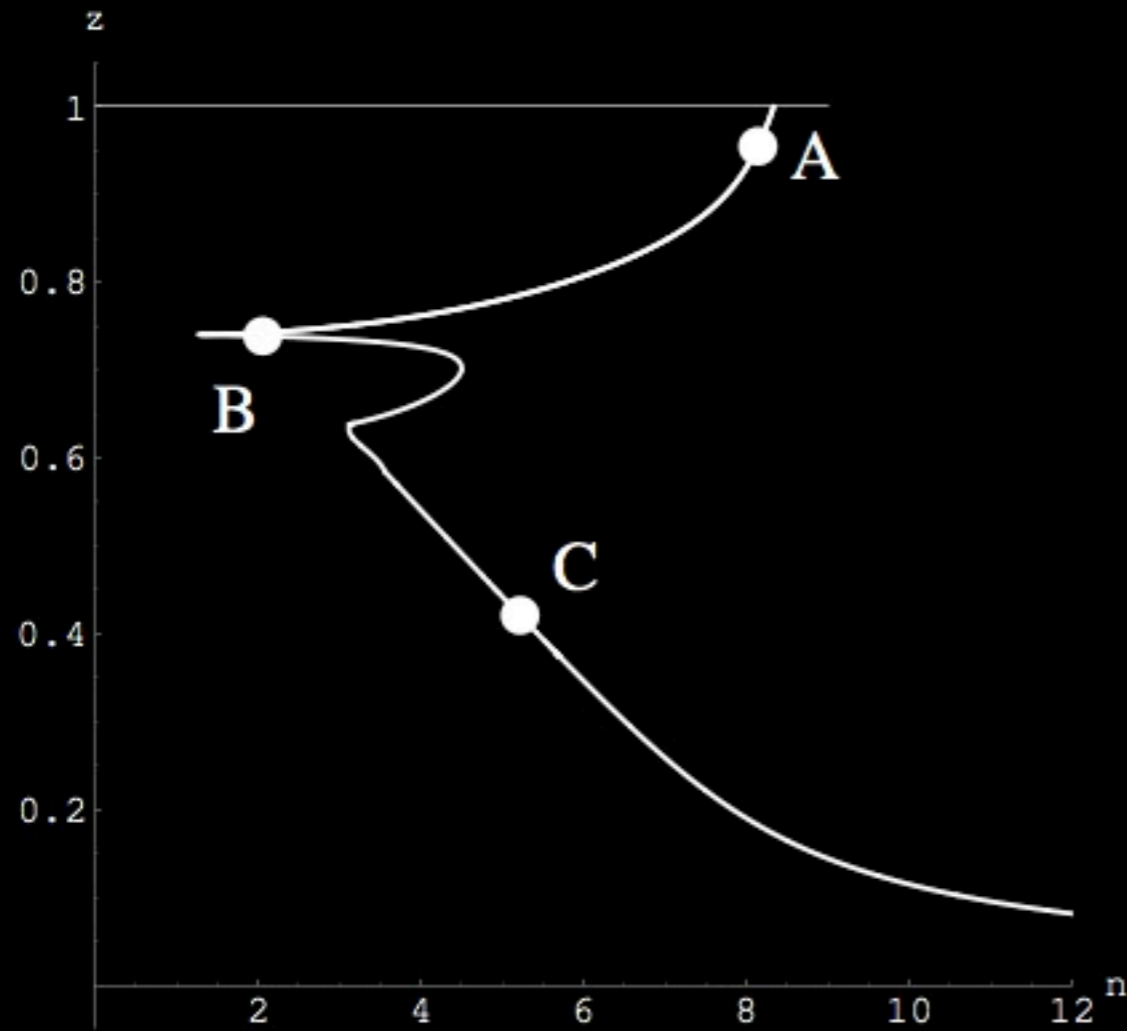


data from Gilles Charvin (LPS-ENS) $\sigma = \frac{n}{n_0}$

Numerical simulations

slenderness ratio: $\frac{L}{2R} = 170$

constant tension: $t = \frac{TL^2}{4\pi^2 K_0} = 6$



(based on Swigon+Coleman model for contact in Kirchhoff rods)

S. Neukirch, "Extracting DNA ...", Phys. Rev. Lett. 93 (2004)

Orders of magnitude

Buckling threshold
for a clamped beam

$$T = (2\pi)^2 \frac{K_0}{L^2}$$



$$\Rightarrow t = \frac{TL^2}{(2\pi)^2 K_0} = 1$$

DNA in tweezers experiments

$$L \simeq 1 \mu\text{m} \quad (\text{few kbp})$$

$$K_0 \simeq 50 \text{ nm } k_B T$$

$$T \simeq 1 \text{ pN}$$

$$\Rightarrow t \sim 10^4$$

Analytical model for plectonemic DNA



straight tails

DNA

L_p



no
end-
loop

uniform ply

$Z = L - L_p$

Elastic rod with :

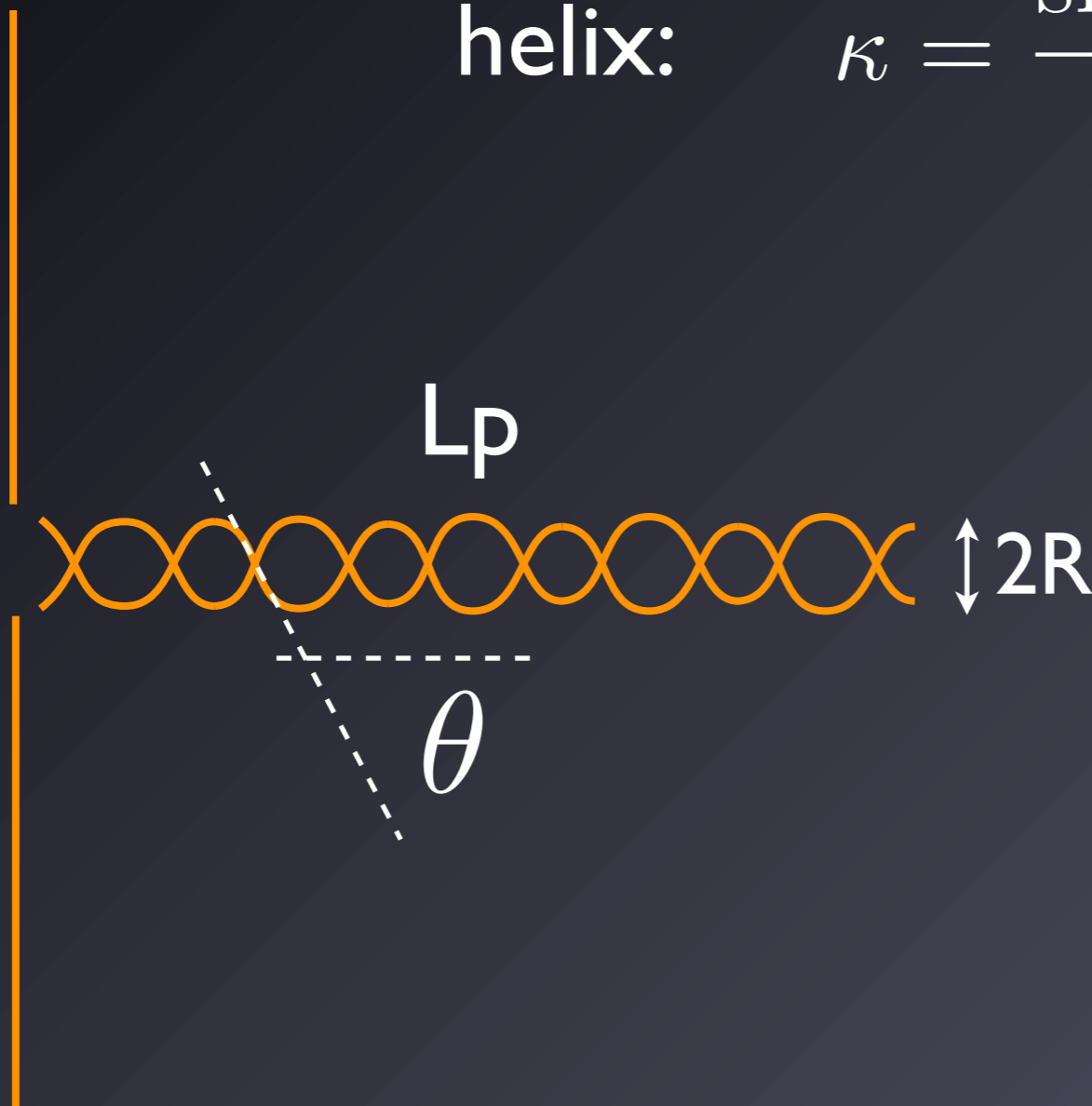
- total length L
- circular cross-section R_0
- bending rigidity K_0
- twist rigidity K_3

N. Clauvelin et al,
Biophysical Journal (2009)

Energy formulation: elastic strain energy

bending: $V = \frac{1}{2} K_0 \kappa^2 L_p$

helix: $\kappa = \frac{\sin^2 \theta}{R}$ (uniform)



Energy formulation: elastic strain energy

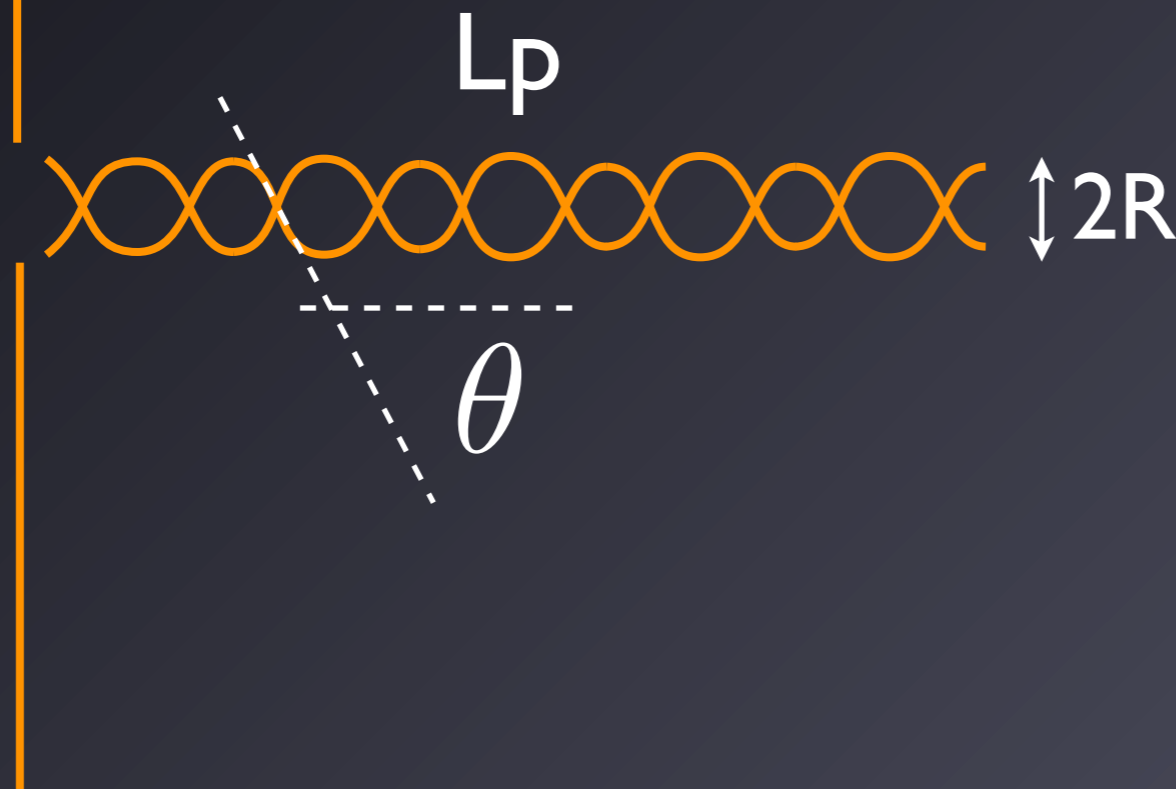


twisting: $V = \frac{1}{2} K_3 \tau^2 L$

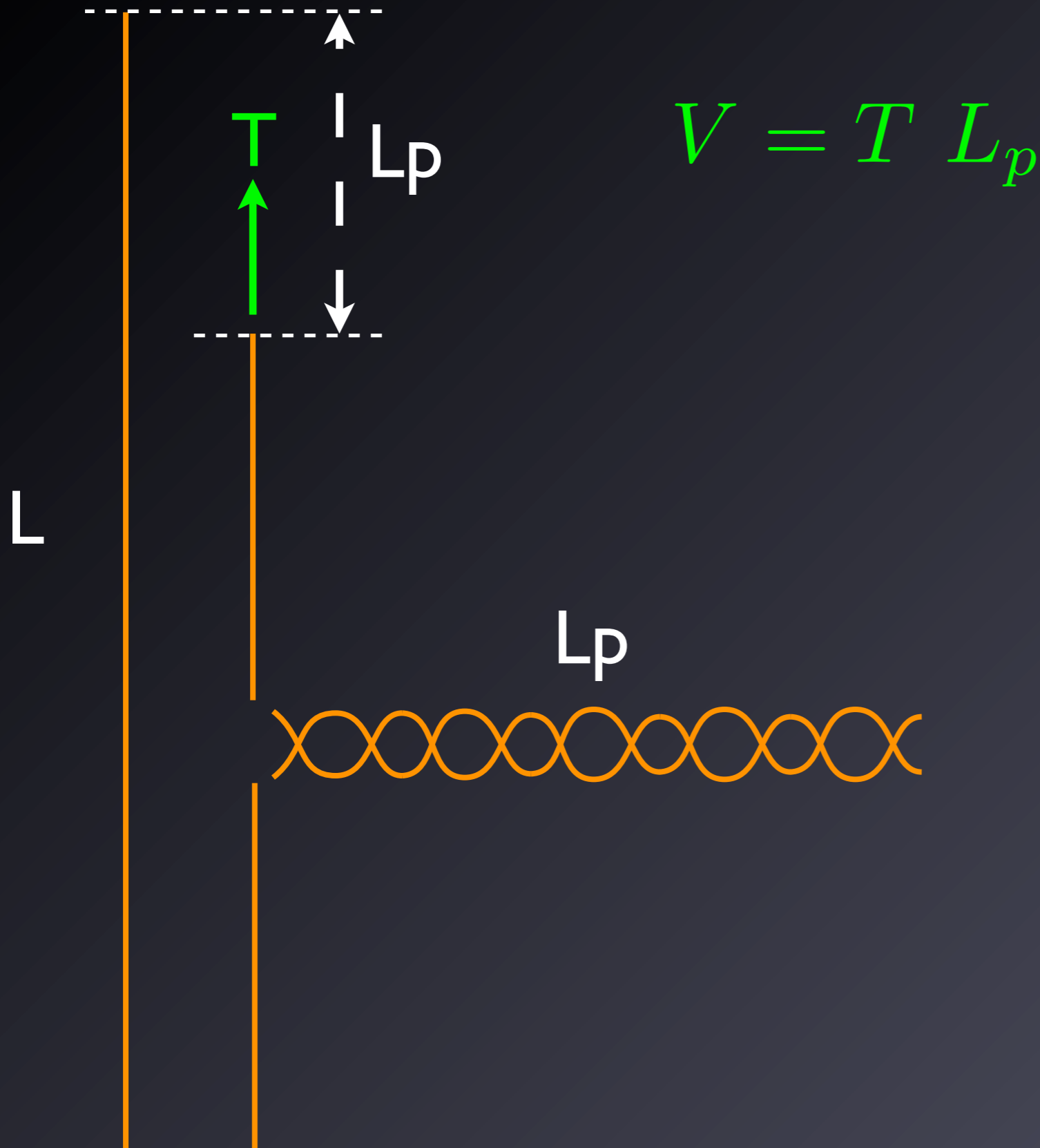
twist τ is uniform along the rod



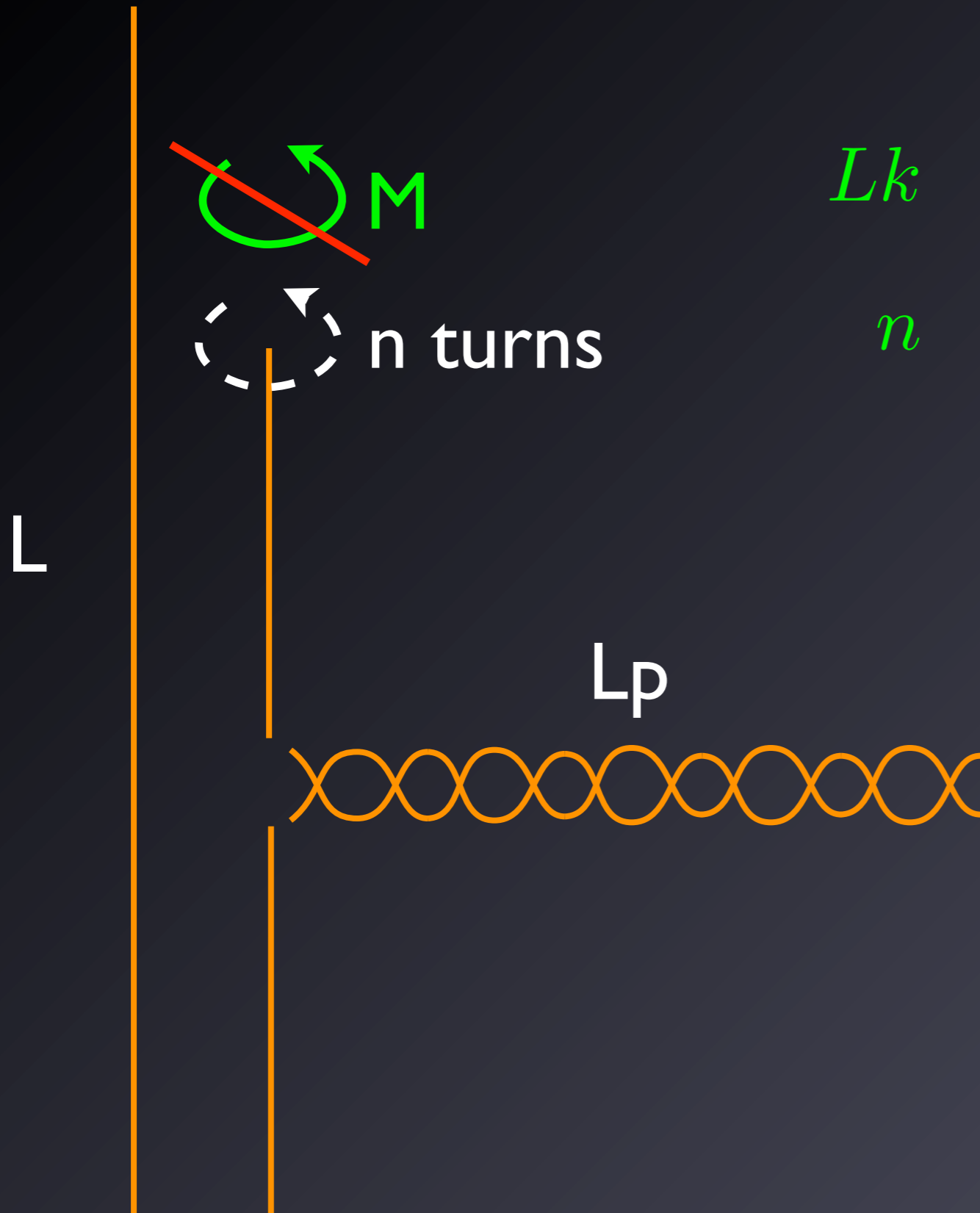
constitutive relation: $M = K_3 \tau$



Energy formulation: work of external loads



Energy formulation: link constraint



$$Lk = Tw + Wr$$

$$n = \frac{1}{2\pi} \left(\tau L + \frac{\sin 2\theta}{2R} L_p \right)$$

Energy formulation: self-interaction

hard-wall (contact)

=> constraint:

$$V = \lambda (R - R_0)$$

long-range:

~ electrostatics

▶ S. Leikin

▶ D. Stigter

▶ Debye-Hückel

▶ G. Manning

▶ ...

$$V = L_p U(\theta, R)$$

Energy formulation: equilibrium

$$V(\theta, R, \tau, L_p) = \frac{1}{2} K_0 \frac{\sin^4 \theta}{R^2} L_p + \frac{1}{2} K_3 \tau^2 L + T L_p + L_p U(\theta, R)$$

$$\text{with constraint } n = Lk = Tw + Wr = \frac{1}{2\pi} \left(\tau L + \frac{\sin 2\theta}{2R} L_p \right)$$

constraint $\Rightarrow L_p = \dots$

$\Rightarrow V = V(\theta, R, \tau)$

Euler-Lagrange equations : $\left(\frac{\partial V}{\partial \theta}, \frac{\partial V}{\partial R}, \frac{\partial V}{\partial \tau} \right) = 0$

Energy formulation: stability

For some T values,
there are two solutions to: $\left(\frac{\partial V}{\partial \theta}, \frac{\partial V}{\partial R}, \frac{\partial V}{\partial \tau} \right) = 0$

=> we compute the Hessian matrix :

$$H = \begin{bmatrix} \partial_{\theta\theta} V & \partial_{\theta R} V & \partial_{\theta\tau} V \\ \partial_{R\theta} V & \partial_{RR} V & \partial_{R\tau} V \\ \partial_{\tau\theta} V & \partial_{\tau R} V & \partial_{\tau\tau} V \end{bmatrix}$$

and we focus on the stable solution.

Energy formulation: equilibrium

$$\frac{\partial V}{\partial \theta} = 2K_0 \frac{\cos \theta \sin^3 \theta}{R^2} + \frac{\partial U}{\partial \theta} - \frac{2}{\tan 2\theta} \left(\frac{K_0 \sin^4 \theta}{2 R^2} + T + U(R, \theta) \right) = 0$$

$$\frac{\partial V}{\partial R} = T - \frac{K_0}{2 R^2} \sin^4 \theta + R \frac{\partial U}{\partial R} + U(R, \theta) = 0$$

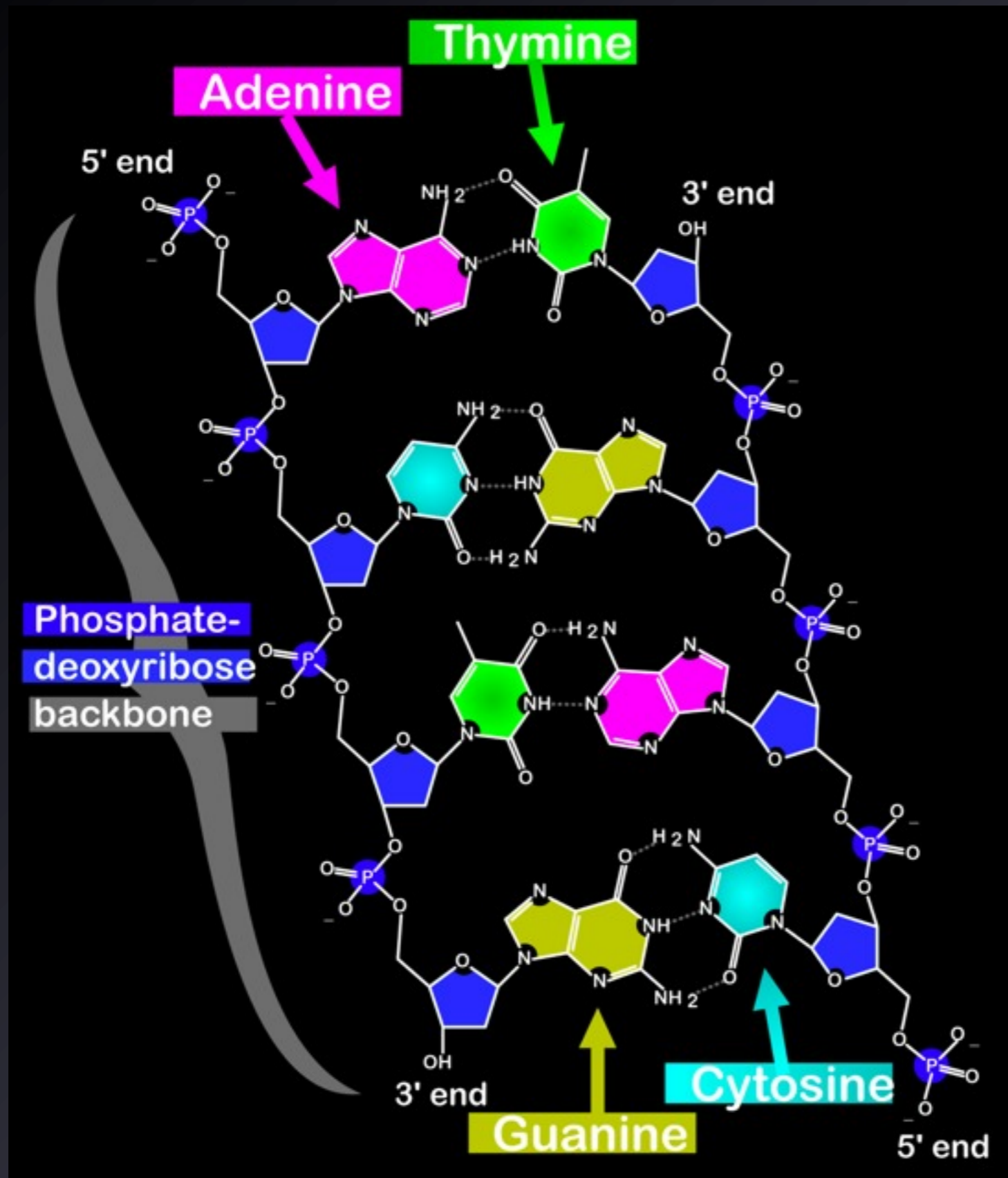
$$\frac{\partial V}{\partial \tau} = K_3 \tau - \frac{2 R}{\sin 2\theta} \left(\frac{K_0 \sin^4 \theta}{2 R^2} + T + U(R, \theta) \right) = 0$$

Once $U(R, \theta)$ is given,

3 equations for

3 unknowns (θ, R, M) ($M = K_3 \tau$)

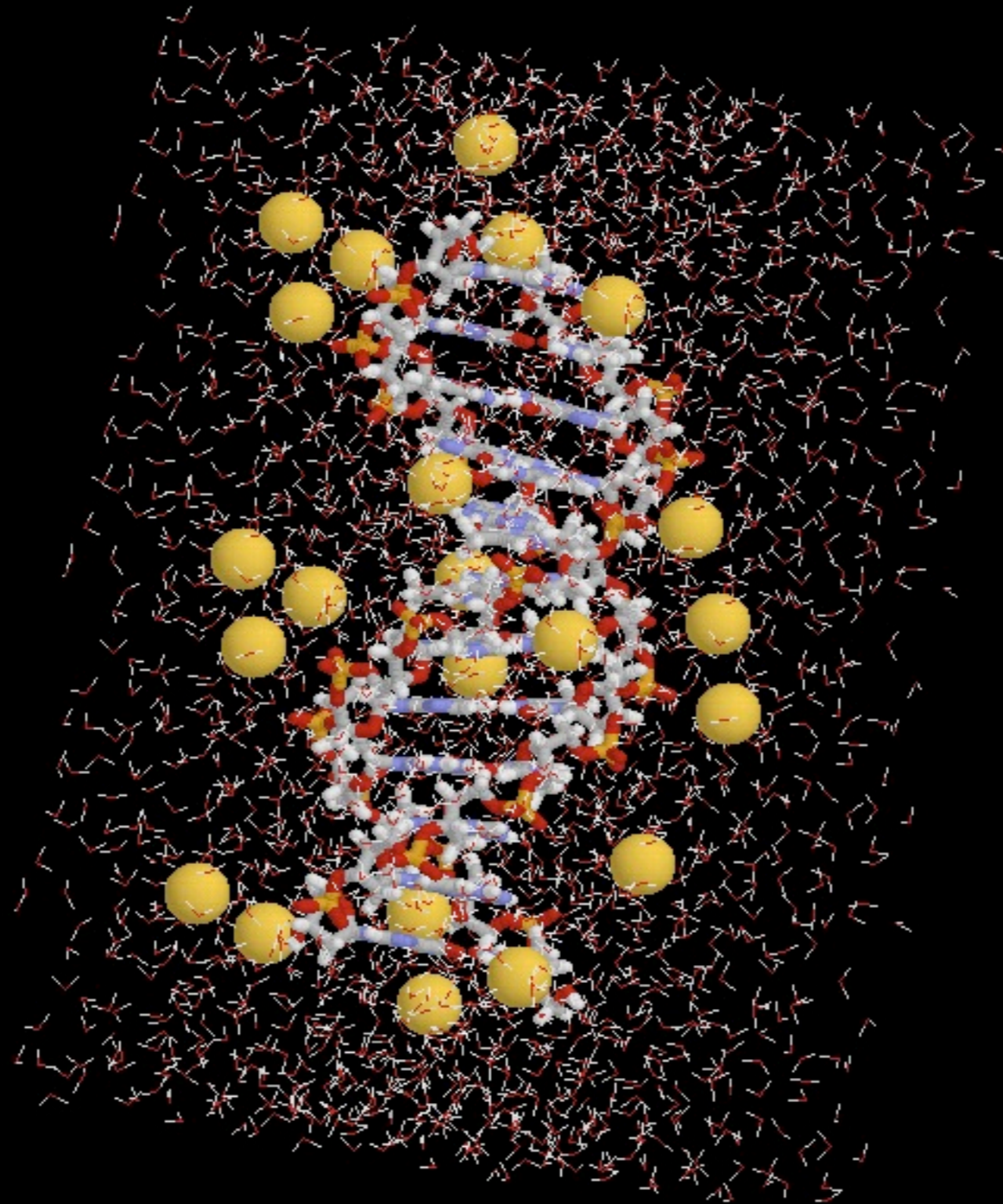
DNA electrostatics



$2 e^-$ per base-pair
 $\Leftrightarrow 1 e^- / 0.17 \text{ nm}$

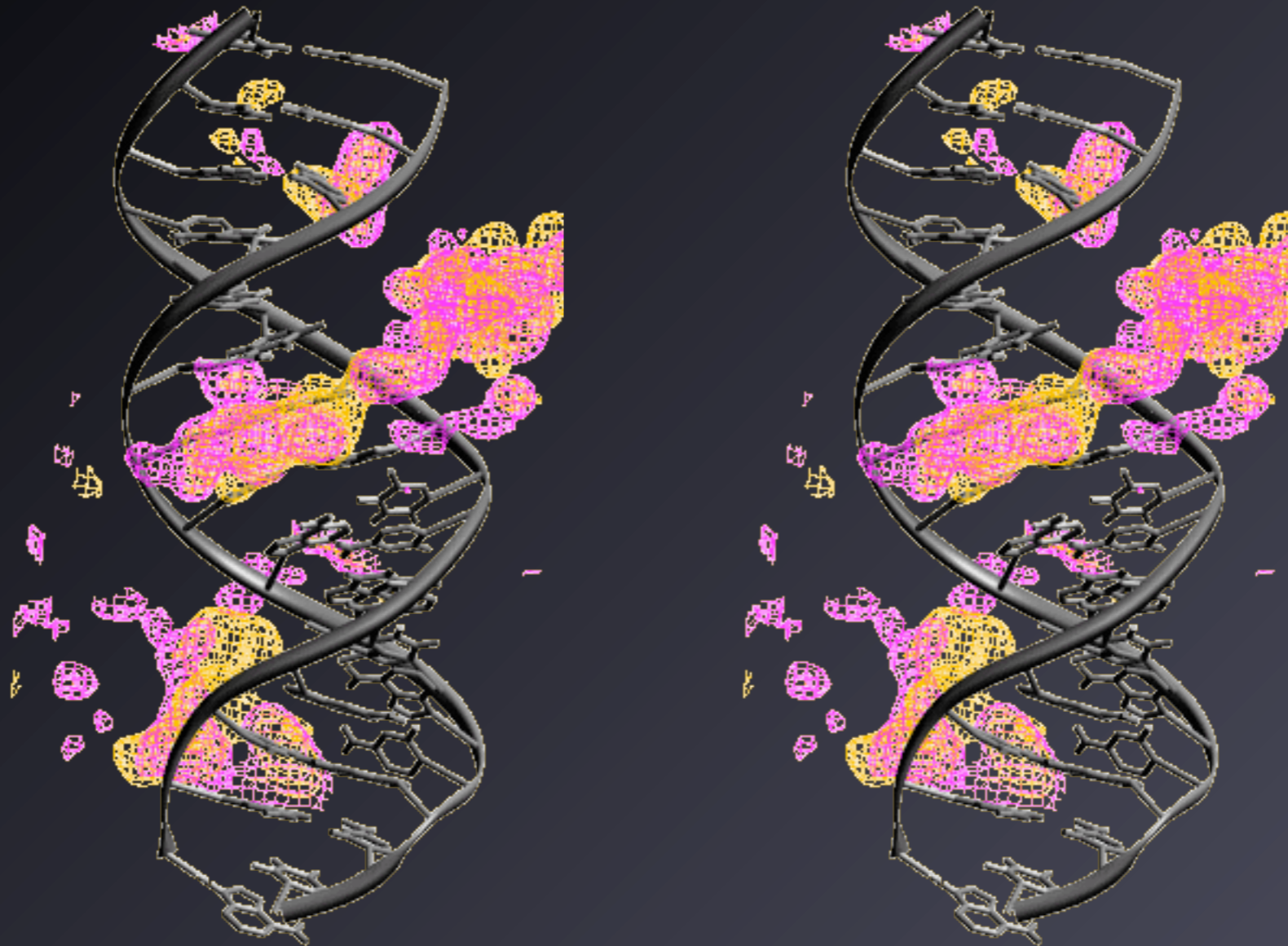
Wikipedia

DNA electrostatics



Alexander
MacKerell
www.psc.edu

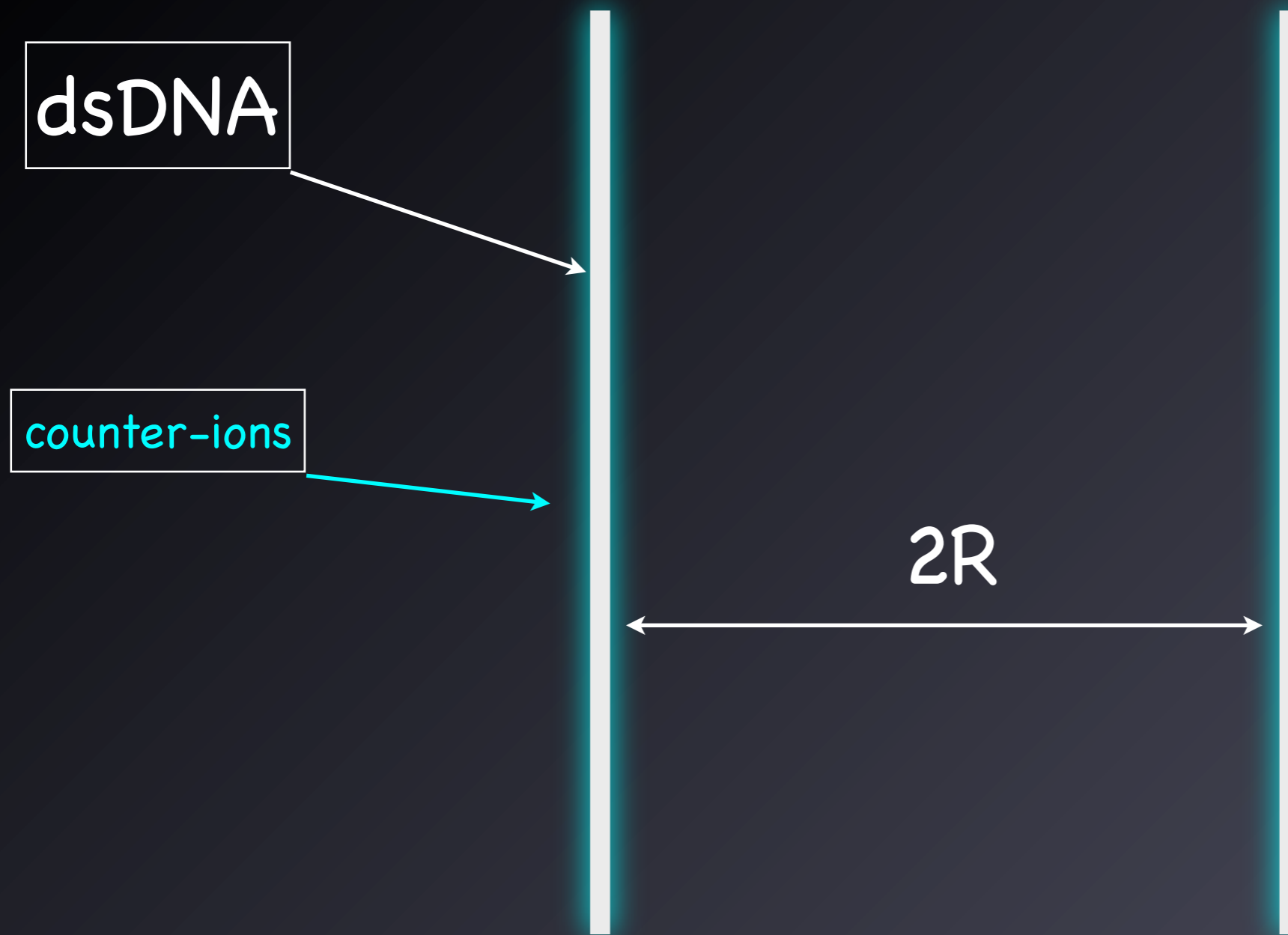
DNA electrostatics



Alexandre Bonvin

www.nmr.chem.uu.nl

DNA electrostatics : Poisson-Boltzmann



effective charge (10mM): $\nu = 1.38/L_B \quad (m^{-1})$

$$L_B = \frac{e^2}{4\pi\epsilon_0\epsilon_r kT}$$

DNA electrostatics : Poisson-Boltzmann

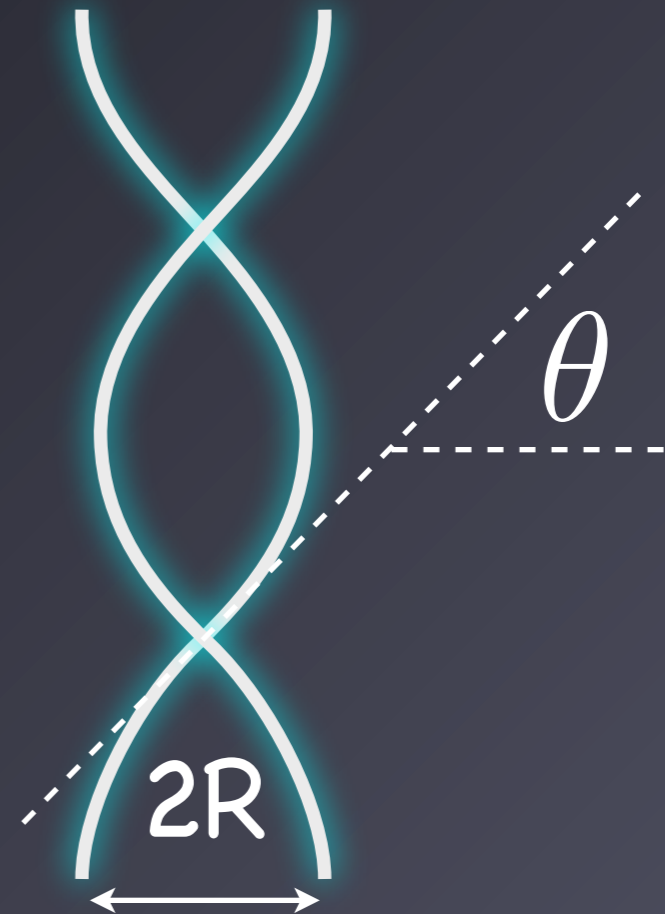
$$U(R, \theta) = \frac{1}{2} kT \nu^2 L_B \sqrt{\frac{\pi}{\kappa R}} e^{-2\kappa R} \cdot \phi(\theta) \quad (\text{per unit length})$$

κ : Debye

$$L_B = \frac{e^2}{4\pi\epsilon_0\epsilon_r kT}$$



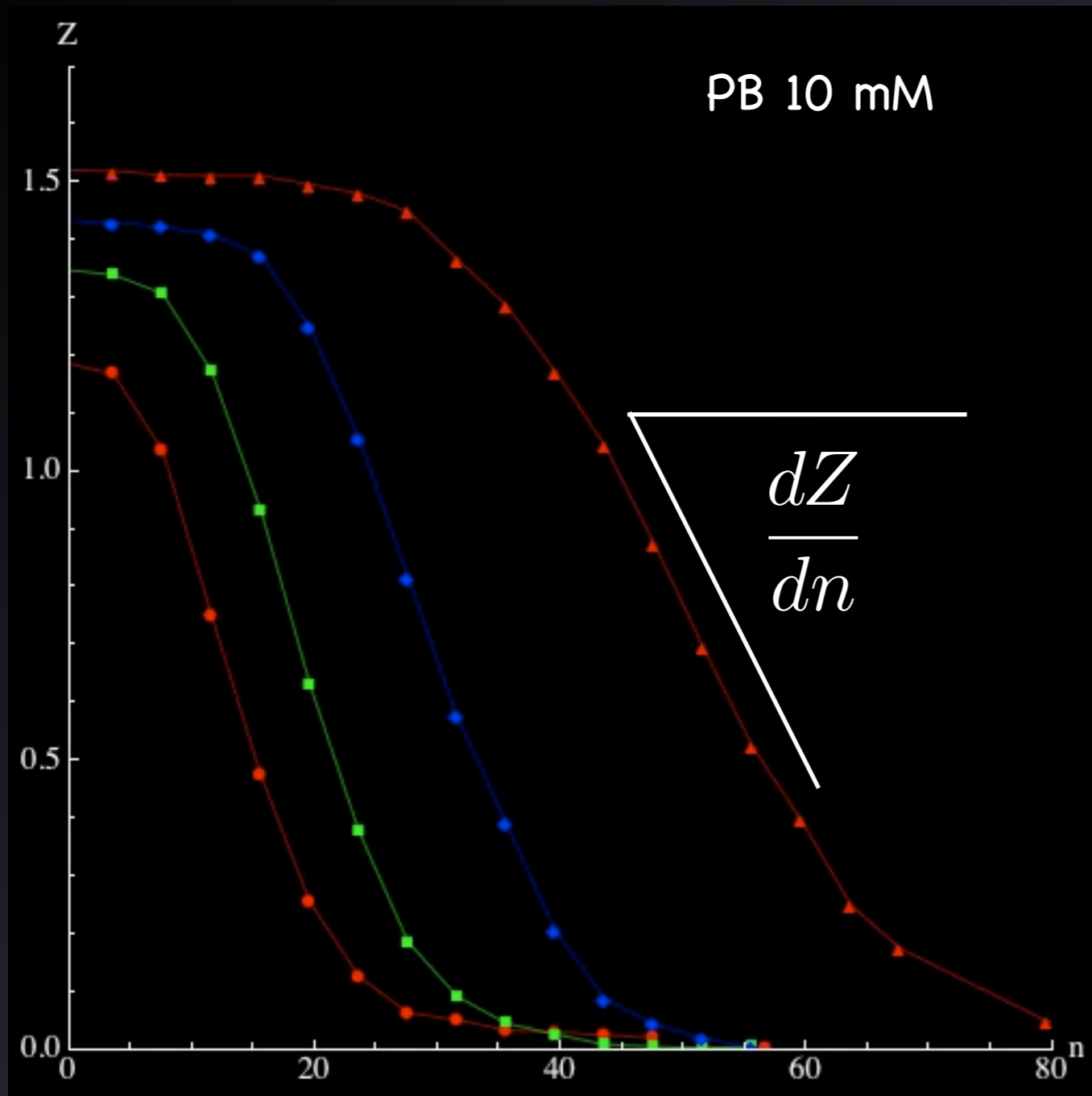
$$\phi(\theta) = 1$$



$$\phi(\theta) = 1 + 0.83 \tan^2 \theta + 0.86 \tan^4 \theta$$

J. Ubbink, T. Odijk, Biophysical Journal (1999)

Results : comparison with experimental data

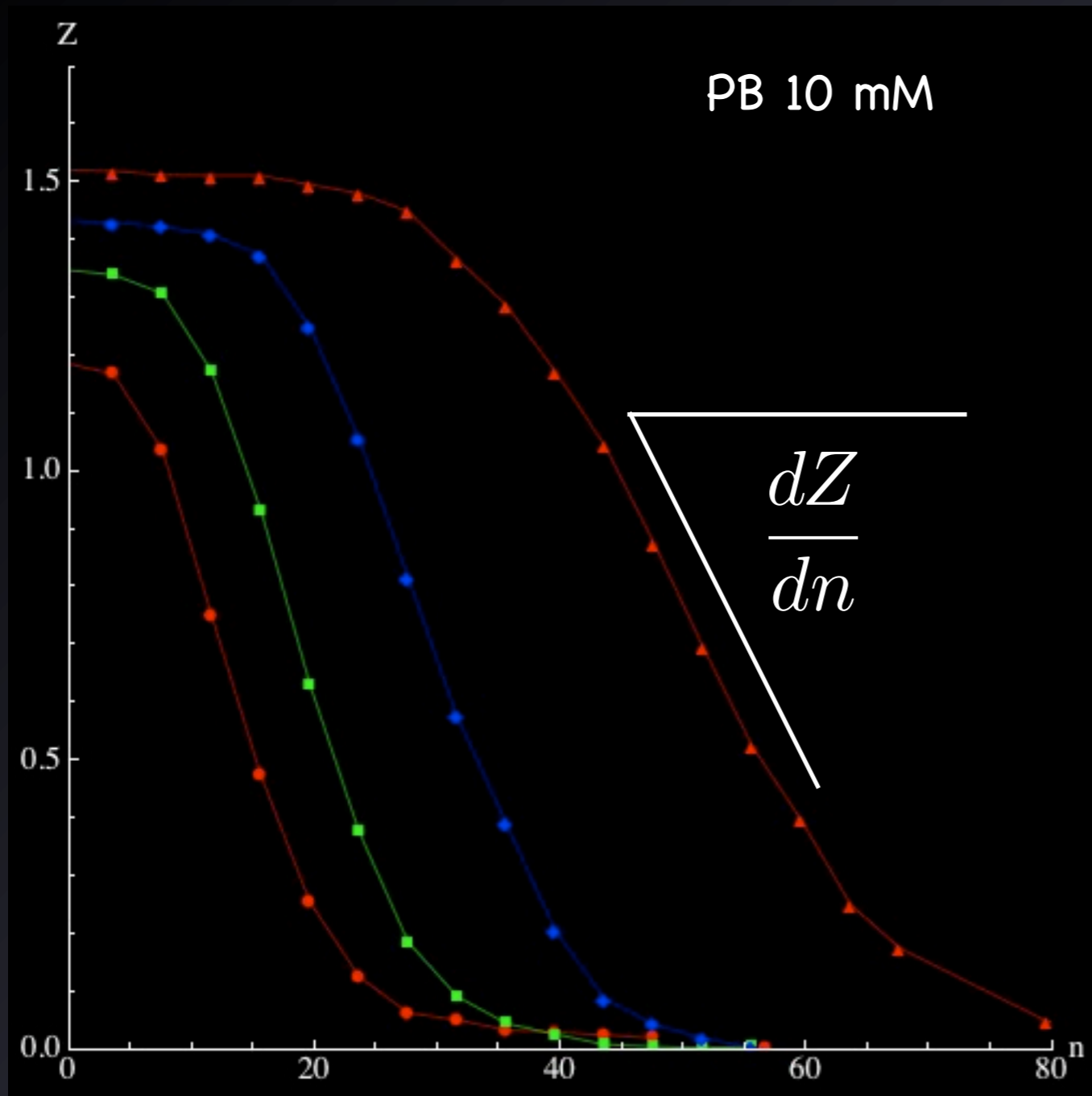


$n = Lk = Tw + Wr$
with Wr linear of z

$$\Rightarrow \frac{dZ}{dn} = \frac{4\pi R}{\sin 2\theta}$$

data from Gilles Charvin (ENS-Paris)

Results : comparison with experimental data



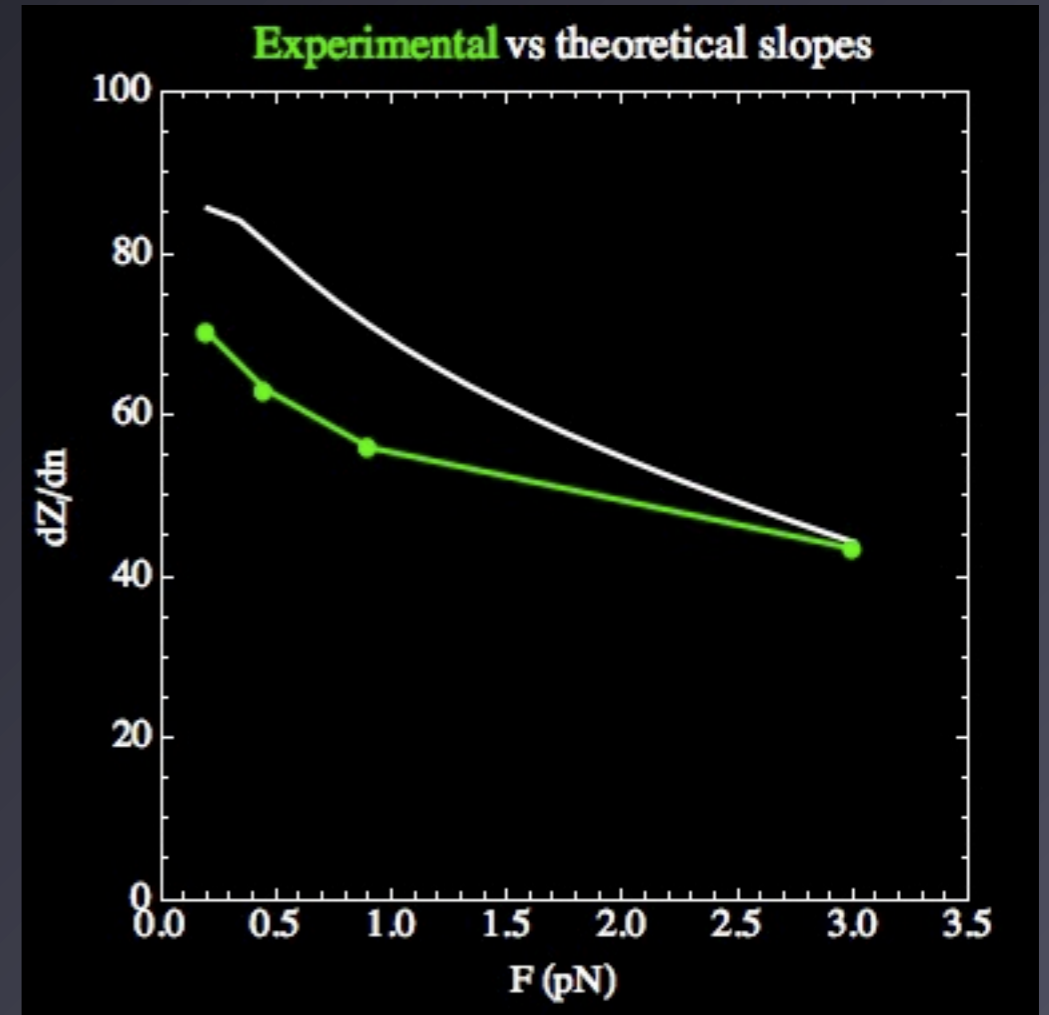
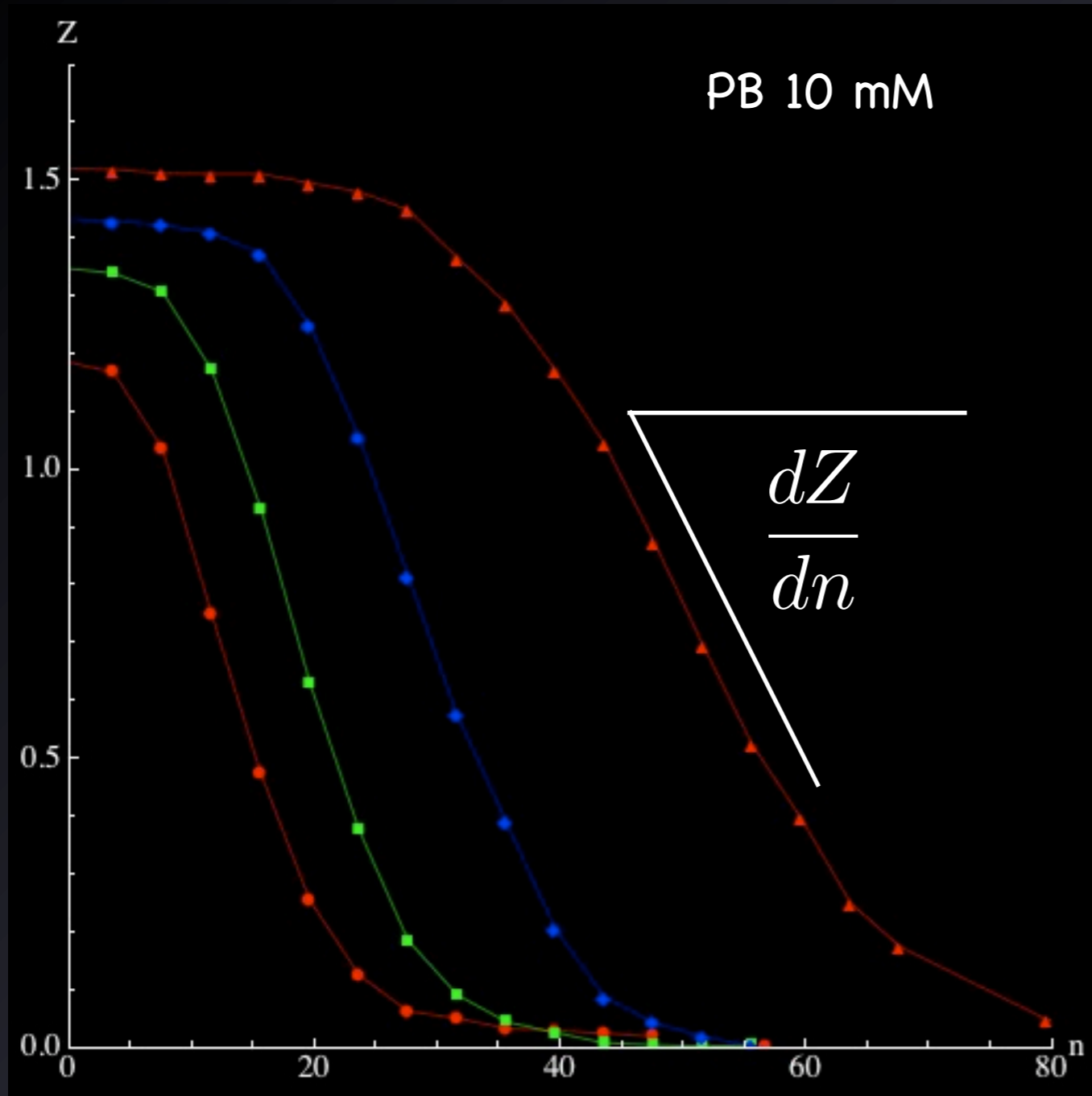
$n = Lk = Tw + Wr$
with Wr linear of z

$$\Rightarrow \frac{dZ}{dn} = \frac{4\pi R}{\sin 2\theta} \rho_{WLC}$$

$$\rho_{WLC} = \frac{Z(n=0)}{L}$$

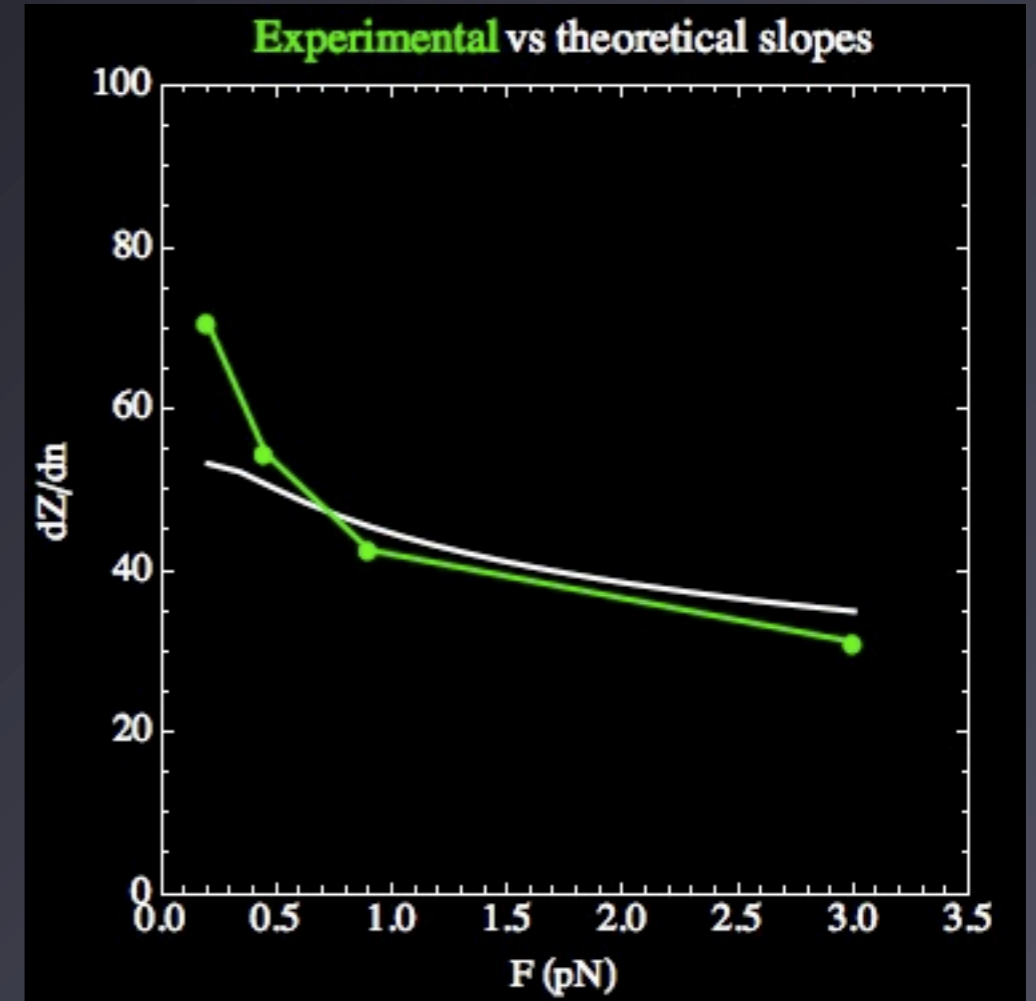
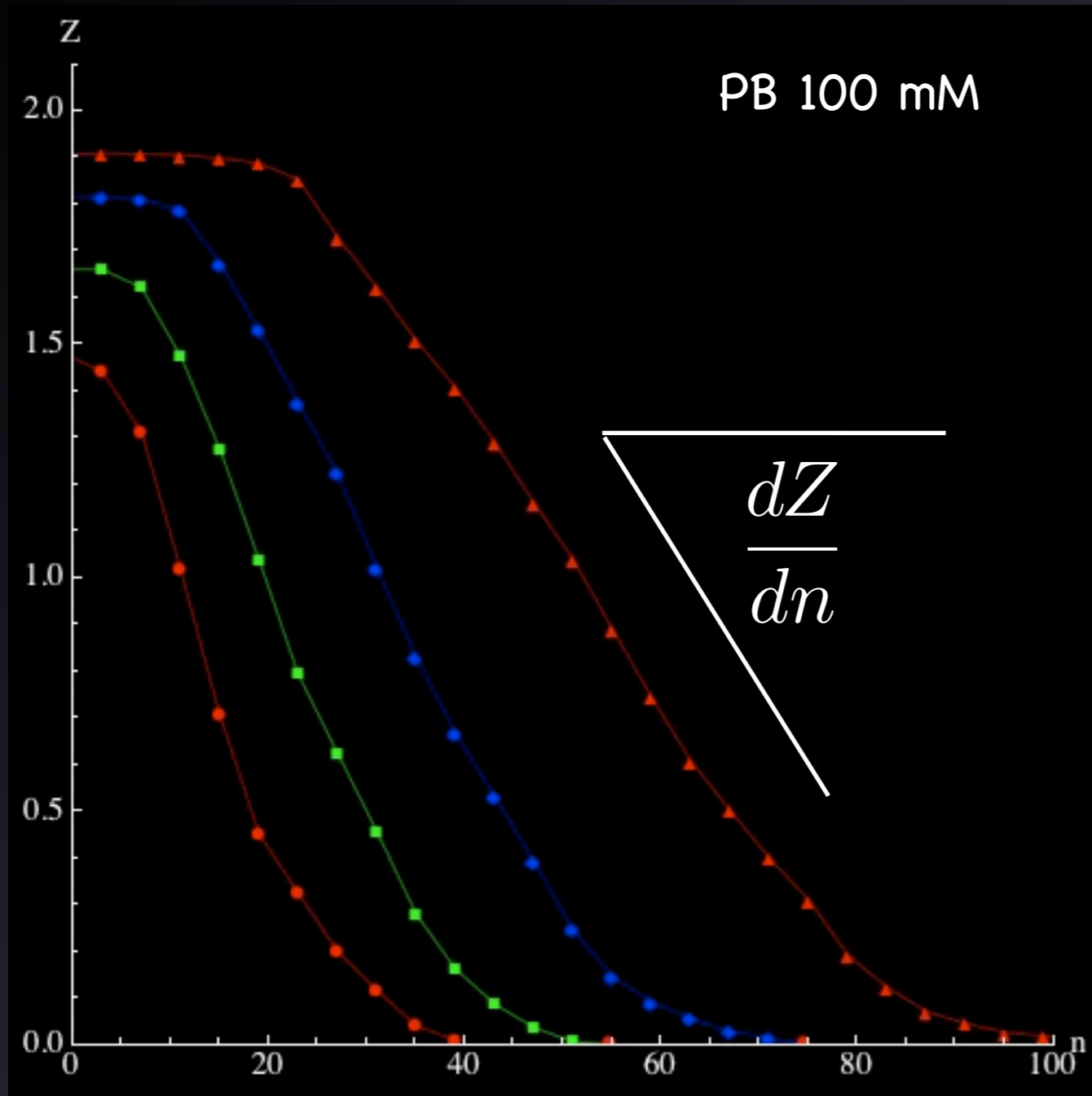
data from Gilles Charvin (ENS-Paris)

Results : comparison with experimental data



data from Gilles Charvin (ENS-Paris)

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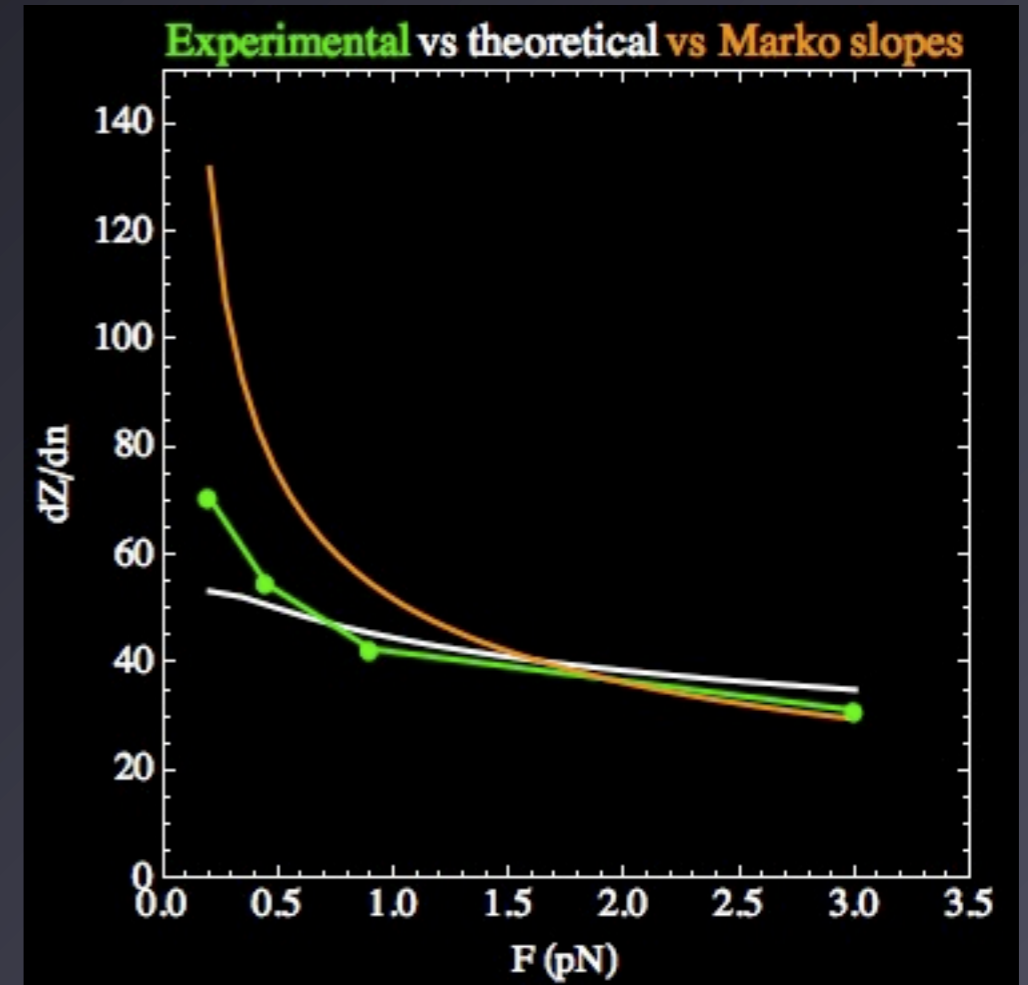
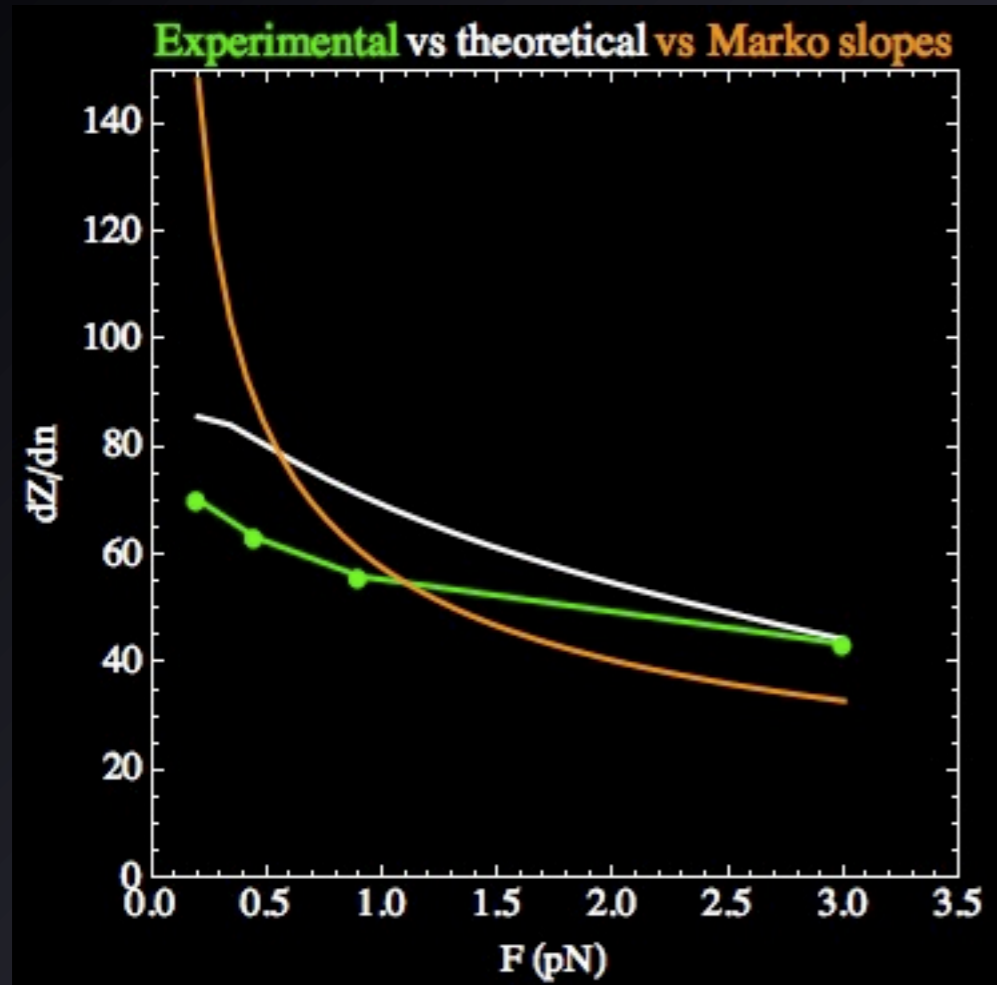


data from Gilles Charvin (ENS-Paris)

Results : comparison with Marko model

PB 10 mM

PB 100 mM



J. Marko, "Torque and dynamics of linking number ...", Phys. Rev. E. (2007)

Remarks

▶ Supercoiling radius R is always $> 1\text{nm}$ (no DNA-DNA contact)

T (pN)	0.2	0.45	0.9	3
R (nm)	3.8	3.3	3.0	2.3

PB 100 mM

▶ Benchmark for DNA-DNA potentials:

1. propose a potential $U(\theta, R)$
2. compute theoretical slopes
3. compare with experiments

Conclusion

- ▶ Analytical model for plectonemic DNA
- ▶ Long-range DNA-DNA interaction potential
- ▶ Reproduces experimental curves (10–100 mM)
- ▶ Could serve as a benchmark for DNA-DNA potentials

- ▶ Thermal fluctuations